

Probing the 3-D matter distribution at $z \sim 2$ with QSO multiple lines of sight†

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Abstract. We investigate the 3-D matter distribution at $z \sim 2$ with high resolution ($R \sim 40,000$) spectra of QSO pairs and groups obtained with the UVES spectrograph at ESO VLT. Our sample is unique for the number density of objects and the variety of separations, between ~ 0.5 and 7 proper Mpc†. We compute the real space cross-correlation function of the Lyman- α forest transmitted fluxes. There is a significant clustering signal up to ~ 2 proper Mpc, which is still present when absorption lines with high column density ($\log N \geq 13.8$) are excluded.

1. Introduction

The Ly- α forest absorbers detected at $z > 1.5$ in high resolution QSO spectra outnumber any other population of objects observable from the ground at those redshifts. They originate, as shown by hydro-simulations (e.g. Cen *et al.* 1994), in density fluctuations of the intergalactic medium (IGM), i.e. in the web of filamentary structures connecting the densest peaks, and are reliable tracers of the baryon density field as well as of the underlying dark matter distribution.

From the study of absorption spectra along single lines of sight (LOSs) to distant QSOs it has been possible to determine the shape and amplitude of the power spectrum of the spatial distribution of dark matter (e.g. Kim *et al.* 2004, Viel, this conference), and gain important information on the baryon density of the Universe (Rauch *et al.* 1997) and on the physical state of the IGM (Schaye *et al.* 2000).

Our present analysis introduces the use of adjacent LOSs to probe the actual 3-D distribution of matter in the Universe and provide estimates of the size/correlation of the absorbing structures. Hydrodynamical (Charlton *et al.* 1997) and analytical (Viel *et al.* 2002) simulations have shown the advantages in using multiple LOSs with respect to single ones, in particular, to decrease the effects of inaccuracies in the continuum fitting.

† Based on observations collected at the European Southern Observatory, Chile
† We adopt a Λ CDM cosmology with $\Omega_m = 0.3$, $\Omega_\Lambda = 0.7$, and $h = 0.72$.

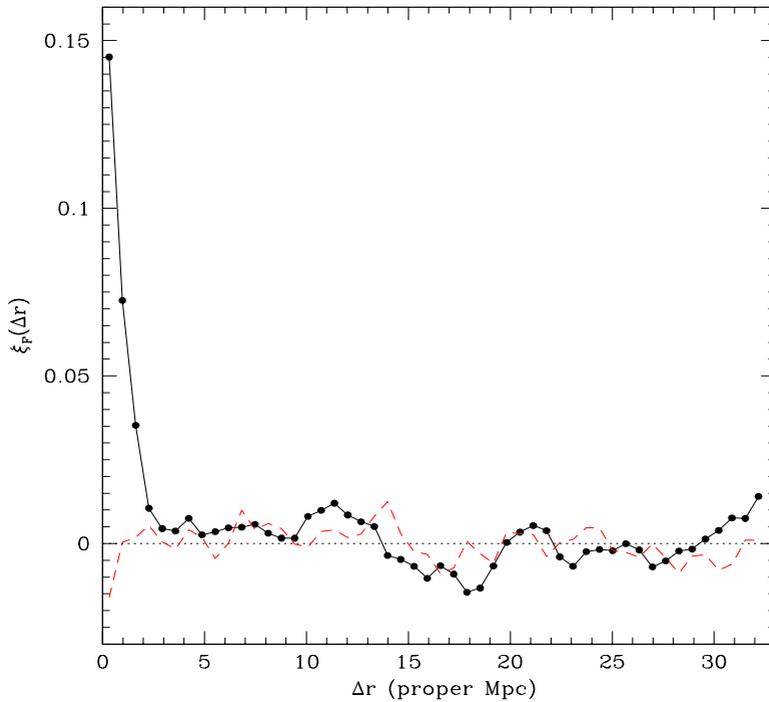


Figure 1. Correlation function of the transmitted flux in real space for our sample of pairs and groups of QSOs. The dashed line is the result obtained for the uncorrelated control sample (see text).

2. Our sample of QSOs

Two major breakthroughs have dramatically improved the exploitation of the potential offered by QSO multiple LOSs: the 2dF QSO Redshift Survey (2QZ, Croom *et al.* 2004), whose complete spectroscopic catalogue contains more than $\sim 23,000$ QSOs, approximately 50 times more than the previous largest QSO survey to a similar depth ($B < 21$) and the UVES spectrograph at the ESO VLT telescope which has a remarkable efficiency especially in the extreme UV.

We have searched the 2dF and other QSO databases for the best groups with apparent magnitude $B \leq 20$ and $z > 1.8$ and carried on a great observational effort to collect UVES spectra of the selected QSOs. At the moment, we have observed one QSO pair, one triplet, and one sextet of QSOs with a resolution $R \sim 40,000$ and signal-to-noise ratio in the Lyman- α forest larger than 3. With the addition of two more UVES QSO pairs from our archive, we total 20 different baselines with proper spatial separations between ~ 0.5 and 7 Mpc, a unique sample both for the number density and the variety of LOS separations investigated.

3. The transmitted flux cross-correlation function

We select in each normalised spectrum the region between the Lyman- β emission (or the minimum observed wavelength, when the Lyman- β falls outside the spectrum) and 5000 km s $^{-1}$ from the Lyman- α emission (to avoid proximity effect) and compute the cross-correlation function of the transmitted flux between adjacent LOSs. The following

generalisation of the formula for the density field (e.g. Peebles 1980) has been adopted:

$$\xi_F(\Delta r) = \frac{\langle (F(x) - \langle F \rangle)(F_0(x + \Delta r) - \langle F_0 \rangle) \rangle}{\sqrt{\langle (F(x) - \langle F \rangle)^2 \rangle \langle (F_0(x + \Delta r) - \langle F_0 \rangle)^2 \rangle}} \quad (3.1)$$

where F and F_0 are the transmitted fluxes in two adjacent LOSs and Δr is the proper spatial separation between two pixels in different LOSs. The cross-correlation function is estimated over all the considered pixels of all the QSO pairs in the sample. To check for the effects of systematics we also computed the same cross-correlation function on a control sample, obtained by substituting in turn one of the QSO spectra in each baseline with an ‘uncorrelated’ one, roughly at the same redshift, but far away on the sky, picked up in our database of high-resolution, high-S/N QSO spectra. The result is shown in Fig. 1. A correlation signal is present up to proper separations of ~ 2 Mpc indicating a large coherence length of the absorbers. This length is consistent with the size of Lyman- α absorbers we found in D’Odorico *et al.* (1998) with the line analysis of coincidences and anti-coincidences.

In order to verify the relative contribution of the stronger and weaker absorption lines to the clustering signal we have selected the pixels with $F > 0.2$ corresponding to $\log N(\text{HI}) < 13.7 - 13.8$ and recomputed $\xi_F(\Delta r)$. Although significantly decreased (by a factor ~ 2) a clustering signal is still clearly present. This is an indication that also far from the most over-dense regions, in the ‘true’ IGM, matter is still clustered.

A comparison with hydrodynamical simulations (see for a description, Viel *et al.* 2004) shows a substantial agreement between the predicted and observed signal.

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