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Towards fast and accurate predictions of radio frequency power deposition and current profile via data-driven modelling: applications to lower hybrid current drive

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Three machine learning techniques (multilayer perceptron, random forest and Gaussian process) provide fast surrogate models for lower hybrid current drive (LHCD) simulations. A single GENRAY/CQL3D simulation without radial diffusion of fast electrons requires several minutes of wall-clock time to complete, which is acceptable for many purposes, but too slow for integrated modelling and real-time control applications. The machine learning models use a database of more than 16 000 GENRAY/CQL3D simulations for training, validation and testing. Latin hypercube sampling methods ensure that the database covers the range of nine input parameters (n_{e0} , T_{e0} , I_p , B_t , R_0 , n_{\parallel} , $Z_{\rm eff}$, $V_{\rm loop}$ and $P_{\rm LHCD}$) with sufficient density in all regions of parameter space. The surrogate models reduce the inference time from minutes to \sim ms with high accuracy across the input parameter space.

Key words: plasma simulation, plasma waves, plasma heating

1. Introduction

The 'advanced tokamak' (AT) concept (Kikuchi 1993; Goldston 1996) is a leading candidate for many steady-state fusion power plant designs (Galambos *et al.* 1995; Sips *et al.* 2005; Najmabadi *et al.* 2006). An AT makes use of the pressure-gradient-driven bootstrap current to sustain a majority of the required plasma current, augmented by non-inductive current drive actuators (Houlberg & Attenberger 1994). These auxiliary actuators may consist of neutral beams and/or radio frequency (RF) systems such as lower hybrid current drive (LHCD) and high harmonic fast wave (HHFW). RF actuators are well

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suited for use in fusion power plants. They do not require line-of-sight access through the breeding blanket and neutron shield surrounding the fusion plasma, and the technology readiness level (TRL) of the RF sources and power delivery systems is high. Control of AT operating scenarios requires detailed information on the RF heating and current drive profiles to maintain magnetohydrodynamic (MHD) stable overall current and pressure profiles (Moreau & Voitsekhovitch 1999; Luce 2011). Moreover, additional current drive is often required at specific radial locations (on axis, off axis, rational q surfaces) to augment the bootstrap current where it is insufficient, or to suppress instabilities such as tearing modes or kink modes.

Consequently, predicting RF wave heating and current drive is critical for present-day fusion experiments and to design and construct a fusion power plant. In fact, the commercialisation of fusion power requires an accurate accounting of plant energy in order to minimise the auxiliary system power usage and to optimise the plant operation cost, which requires advancement in high-fidelity modelling of RF actuators. In addition, accurate predictions of the current drive and heating profiles are needed in integrated modelling to assess the performance and stability of the plasma operating scenario.

An AT plasma is a complex integrated physics system with bi-directional coupling between heat/particle transport, temperature/density profiles, current profile, MHD stability, heat/particle exhaust and fusion reaction rate (Hayashi 2010; Moreau *et al.* 2013; Meneghini *et al.* 2015; Poli 2018). Few of these quantities are directly controllable through external actuators. Control of the RF heating and current drive profiles, by varying the launched power and wavenumber of the system, is one of the few direct control knobs available on a tokamak. Having the capability to determine where the RF system is depositing power/current at the present moment (and how that may change with other plasma conditions) is crucial for the reliable operation of an AT. Although RF simulation tools are sophisticated and accurate, the computational resources required are considerable for high-fidelity simulations.

In this work, we propose using machine learning (ML) models to accurately predict RF power absorption and current density profiles and accelerate the time required for simulations, which will enable widespread use in parametric scenario scoping studies, time-dependent modelling and real-time control. We show that our models are able to quickly and accurately reproduce the results of GENRAY/CQL3D (Harvey & McCoy 1992; Smirnov & Harvey 1995), with inference times on the order of milliseconds as compared with minutes for the ground-truth simulations.

The remainder of this paper is structured as follows. Section 2 reviews the state of surrogate models in fusion research and the physics we reproduce with our surrogate models. Section 3 covers the design and methods used in our research. Section 4 presents the results of our surrogate models and § 5 covers the paper conclusions and future work.

2. Background and previous work

In this section we review the relevant theoretical background on ray tracing Fokker–Plack (FP) modelling of heating and current drive, which is the basis for the input data of our designed ML models. We also review relevant works in the literature that use ML to accelerate computations and we provide a landscape of ML methods in the areas of plasma physics and fusion.

2.1. Ray tracing/FP modelling of heating and current drive

Several approaches are possible to simulate the RF current drive profile in tokamak plasmas. Analytical and reduced numerical (e.g. adjoint method) models provide a reasonable estimate of the total driven current but lack key physics necessary to

determine the current profile shape that can only be realised by solving the FP equation. Moving upward in complexity to ray tracing/FP modelling includes the necessary wave propagation and quasi-linear Landau damping physics for accurate current profile prediction in several frequency ranges (lower hybrid, electron cyclotron). Ray tracing makes use of the simplifications of geometrical optics and reduces the problem from solving Maxwell's equations in toroidal geometry (the 'fullwave' method) to the following set of ray equations:

$$\frac{\mathrm{d}x}{\mathrm{d}\tau} = \frac{\partial D_0}{\partial k},\tag{2.1a}$$

$$\frac{\mathrm{d}\omega}{\mathrm{d}\tau} = \frac{\partial D_0}{\partial t},\tag{2.1b}$$

$$\frac{\mathrm{d}\mathbf{k}}{\mathrm{d}\tau} = -\frac{\partial D_0}{\partial x},\tag{2.1c}$$

$$\frac{\mathrm{d}t}{\mathrm{d}\tau} = -\frac{\partial D_0}{\partial \omega},\tag{2.1d}$$

where D_0 is the wave dispersion relation, ω is the wave frequency, k is the wavenumber, τ is the ray step parameter, t is the time coordinate and x is the spatial coordinate. Equation (2.1b) is trivial under the condition that the time variation of D_0 is small on the timescale of the operating frequency f_0 , which is well satisfied for waves in the gigahertz range.

Solving the FP equation is required for both fullwave and ray-tracing approaches to determine power absorption and current drive. The FP equation

$$\frac{\mathrm{d}f}{\mathrm{d}t} = RF + C,\tag{2.2}$$

describes the time evolution of the distribution function, f, in the presence of RF damping, RF, and collisions, C. The RF term

$$RF = \frac{\partial}{\partial \mathbf{v}} D_{QL}(\mathbf{v}) \frac{\partial}{\partial \mathbf{v}} f_{e0}, \tag{2.3}$$

depends on the distribution function and the quasi-linear diffusion operator, D_{QL} . The wave solvers provide D_{QL} as an input to the FP solver, as shown in figure 1. The FP solver then feeds information on power absorption back to the wave equation solver and the process iterates until convergence. In the ray-tracing approach, the ray trajectories need not be recalculated at each time step of the FP solver, only the power absorption along the ray. A schematic of the workflow between wave equation solvers and FP solvers is shown in figure 1. In this study we use the GENRAY (Smirnov & Harvey 1995) ray-tracing code coupled to CQL3D (Harvey & McCoy 1992; Petrov & Harvey 2016) for solving the FP equation, although the approach may be expanded to include other RF simulation tools such as TORIC (Brambilla 1996) for fullwave simulations.

Both GENRAY and CQL3D require information about the plasma density, temperature and ion composition, as well as the magnetic equilibrium as inputs. In addition, CQL3D requires the toroidal DC electric field, whereas GENRAY requires the wave parameters (power, n_{\parallel} spectrum, antenna position) as initial conditions. The outputs of GENRAY, namely the ray trajectories and quasi-linear diffusion operator, are of indirect rather than direct interest from the perspective of plasma control and integrated modelling. The outputs of CQL3D (wave power deposition and current profiles) are of direct interest and will be the prediction targets of our ML models.



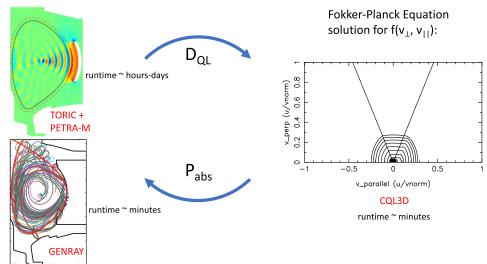


FIGURE 1. Wave equation solvers (TORIC+Petra-M or GENRAY) couple to FP equation solvers (CQL3D) through the quasi-linear diffusion operator, D_{QL} . The FP solver perturbs the distribution function and couples information on quasi-linear RF power absorption, P_{abs} , back to the wave solver. This process iterates until the solutions converge. Rough per-iteration wall clock runtimes for each code are indicated in the figure. TORIC+Petra-M and GENRAY figures reproduced from Shiraiwa *et al.* (2017) and Wallace *et al.* (2010), respectively.

2.2. Using ML to accelerate computations

Modern computational resources have advanced numerical simulation of physical systems to an increasing range of spatial and temporal scales. However, even with high-performance computing techniques, it is a challenge for real-time demand response in feedback control. Emerging data-driven methods have revolutionised how researchers model, predict and eventually control these complex systems in a diverse range of fields, including but not limited to automation, climate, combustion, fluids, high-energy physics, plasma science and plasmonics (Malkiel *et al.* 2018; Duraisamy, Iaccarino & Xiao 2019; Wilkes *et al.* 2020; Yellapantula *et al.* 2020; Abbate, Conlin & Kolemen 2021; Bai & Peng 2021; Hatfield *et al.* 2021; Maschler & Weyrich 2021; Brunton & Kutz 2022). A variety of ML techniques have been used such as deep neural networks (DNNs) to fit simulation data. The resulting surrogate models have been used in larger frameworks. For example, global climate models use ML models to represent kernels drawn from observations (Reichstein *et al.* 2019). A common thread is the use of numerical models and codes to generate training data to produce surrogate models with a low mean squared error (MSE). This is the approach of this paper.

2.3. The landscape of ML in plasma physics and fusion

The application of ML techniques to fusion energy and plasma physics dates to the late 20th century, with initial applications of neural networks to the topics of tokamak plasma control (Bishop *et al.* 1992, 1995*a*; Wróblewski 1997), disruption prediction (Hernandez *et al.* 1996; Wroblewski, Jahns & Leuer 1997), energy confinement prediction (Allen & Bishop 1992) and data analysis (Lister & Schnurrenberger 1991; Bishop *et al.* 1993; Morabito & Versaci 1997), among other topics. These early applications showed the promise of ML techniques in fusion research. Considerable advances in both computing

hardware and ML techniques after the turn of the century resulted in renewed interest in ML applications for fusion in recent years.

ML continues to be an extremely effective tool in some of the early adopted topics, particularly disruption prediction (Rea & Granetz 2018; Rea et al. 2020; Churchill, Tobias & Zhu 2020a; Churchill et al. 2020b). A more recent application of ML to fusion research is in the area of accelerating computationally intensive simulation models. Researchers developed fast surrogate models for plasma transport (Meneghini et al. 2020; van de Plassche et al. 2020) and neutral beam heating/current drive (Boyer, Kaye & Erickson 2019; Morosohk, Boyer & Schuster 2021), used surrogate model frameworks for validation of plasma transport codes (Rodriguez-Fernandez et al. 2018) and also employed ML techniques to reduce computation time for codes such as XGC (Miller et al. 2021). The work of Boyer and Morosohk, in which they create a fast neural-network=based surrogate model for NUBEAM, is particularly relevant to the work presented in this paper. The work presented here represents the first efforts to create a ML-based surrogate model for RF heating and current drive.

3. Design and methods

To reduce the computational cost of ray tracing/FP modelling of heating and current drive, we implement and study three ML models to tackle the regression problem to predict both current profiles and wave power deposition. In the following, we first describe the details of the source data and input/output parameters used to design the ML models. Next, we describe the three ML models we implement in this study, multilayer perceptron (MLP), Gaussian processes and random forest, along with information about their design and hyperparameter optimisation.

3.1. Source data

Training ML models requires a database of simulations to draw training, validation and testing data. In the ML context, 'validation' refers to the process of using a holdout set of data to provide an unbiased evaluation of the model while tuning the hyperparameters. To this end, the π Scope (Shiraiwa *et al.* 2018) workflow engine generates a set of 16 384 GENRAY (Smirnov & Harvey 1995)/CQL3D (Harvey & McCoy 1992) input files and submits five of these simulations at a time to run in parallel on the MIT Engaging cluster at the Massachusetts Green High Performance Computing Center. The wall-clock time required to generate the entire database was approximately two weeks. As each GENRAY/CQL3D simulation completes, π Scope incorporates the relevant input and output quantities into a NETCDF database. Because a single NETCDF file containing all simulation data would be too large for efficient data read/write, π Scope chunks the NETCDF files with 100 simulations per file.

The GENRAY/CQL3D input files contain hundreds of parameters covering both the plasma and the numerical methods used in the simulation. The vast majority of these parameters do not vary in the simulation database, with nine zero-dimensional parameters describing the plasma and lower hybrid system varying across the ranges listed in table 1. The simulation database consists of LHCD simulations for the EAST tokamak (Wu *et al.* 2007). The LHCD wave frequency is 4.6 GHz, and the antenna size/location match that of the LHCD antenna on EAST (Liu *et al.* 2015). The kinetic profiles are analytic and defined by a central electron temperature/density (n_{e0} , T_{e0}) and edge density/temperature (n_{eb} , T_{eb}) according to the equations

$$n_e = (n_{e0} - n_{eb}) \times (1 - \rho^2) + n_{eb}$$
 (3.1)

Parameter	Symbol	Range	Units
Magnetic field on axis	B_0	2.5-3.5	T
Plasma current	I_p	0.4 - 1.2	MA
Parallel refractive index	n_{\parallel}	1.7-2.5	
Central electron density	n_{e0}	$1-5 \times 10^{19}$	m^{-3}
Central electron temperature	T_{e0}	1–5	keV
Plasma effective charge	$Z_{ m eff}$	1.5-2.5	
Plasma major radius	R_0	1.8-1.9	m
Toroidal DC electric field	E_{DC}	-0.0001- 0.001	$\rm Vcm^{-1}$
Lower hybrid wave power	$P_{ m LH}$	0.1-3	MW

TABLE 1. Parameters varied in the GENRAY/CQL3D database.

and

$$T_e = (T_{e0} - T_{eb}) \times (1 - \rho^2) + T_{eb},$$
 (3.2)

where ρ is the normalised minor radius. In the GENRAY/CQL3D database n_{e0} and T_{e0} vary whereas n_{eb} and T_{eb} remain constant at 7.5 \times 10¹⁸ m⁻³ and 0.1 keV, respectively.

The magnetic equilibrium is an analytic model based on the work described in (Guazzotto & Freidberg 2021). The major radius of the plasma, plasma current and toroidal magnetic field vary over typical ranges for the EAST tokamak in this database. However, other equilibrium parameters such as inverse aspect ratio (0.24), triangularity (0.5), elongation (1.8) and magnetic configuration (double null) remain constant for all simulations in the database.

For sampling the parameter space, we chose to use Latin hypercube sampling (LHS) (McKay, Conover & Whiteman 1976). LHS covers a large dimension parameter space with far fewer sample points than a uniform grid and guarantees better uniformity across all columns/rows of the multi-dimensional parameter space than random sampling methods. Given the utility of this sampling method, π Scope now incorporates LHS functionality in parametric scans of input variables.

The parameter ranges in table 1 cover or exceed typical values for each parameter seen on the EAST tokamak, however not all combinations of parameters are relevant. In particular, a significant number of the simulations occur in the 'slide-away' regime where the collisional drag on fast electrons is insufficient to counteract the accelerating potential of the toroidal DC electric field:

$$m_e v v_{\text{coll}}(v) < q E_{\text{crit}},$$
 (3.3)

where m_e is the electron mass, v is the fast electron velocity, $v_{\text{coll}}(v)$ is the fast electron collision frequency, q is the electron charge and E_{crit} is the critical toroidal DC electric field. Using

$$\nu_{\text{coll}}(v) = \frac{q^4 n_e (\ln \Lambda)}{4\pi \epsilon_0 m_e^2 v^3},$$
(3.4)

results in

$$E_{\text{crit}} = \frac{q^3 n_e (\ln \Lambda)}{4\pi \epsilon_0 m_e v^2} \tag{3.5}$$

which is the familiar expression for the Dreicer field (Dreicer 1959). Taking $v = c/n_{\parallel}$ as the upper limit of the quasi-linear plateau for slide-away electrons (as opposed to v = c

for true runaway electrons (Connor & Hastie 1975)) yields

$$E_{\rm crit} = 7.9 \times 10^{-22} n_e n_{\parallel}^2 \tag{3.6}$$

with $E_{\rm crit}$ measured in V m⁻¹. The slide-away regime can be modelled accurately with CQL3D by increasing the resolution and range of the electron energy grid, however this requires additional computational resources and the slide-away regime is generally avoided in LHCD experiments. A threshold of $E_{\rm DC}/E_{\rm crit}>0.9$ delineates the slide-away regime to be filtered from the dataset, resulting in the final dataset used to design the ML models as described in the following.

3.2. ML models

This section describes the three different ML models designed to predict current profiles and power deposition. Each of these ML models carry different characteristics which may affect prediction performance in different ways. Our goal is to evaluate multiple ML models as surrogates for the traditional more expensive type of computation described in the previous section.

3.2.1. MLP

MLP (Bishop *et al.* 1995*b*) is a feed-forward artificial neural network composed of multiple layers of neurons (perceptrons) with activation functions. It typically consists of three types of layers of nodes, an input layer, one or more hidden layers and an output layer. Each node (with exception of the input layer) is a neuron that utilises a nonlinear activation function to learn the right representation of the data mathematically. MLP provides global approximations by mapping a set of input features to output features through a training process using the back-propagation algorithm (Rumelhart, Hinton & Williams 1986). During this process, the parameters consisting of weights and biases are iteratively adjusted and optimised. MLP model has been broadly applied in classification and regression tasks in science (Gardill *et al.* 2018; Yüksel, Soydaner & Bahtiyar 2021). MLP is a supervised learning algorithm that relies on the existence of training data and is suitable for use in both classification and regression problems.

The hyperparameters of the neural network include the batch size, learning rate, the L_2 normalisation factor, the number of hidden layers and the number of neurons in each hidden layer. Considering the large search space of the hyperparameters in MLP, we employ the random search method by placing a reasonable set of range of the hyperparameters. We constrain the MLP to having four hidden layers, where the number of nodes in each layer is searched in the range [20, 1000], the learning rate in [10⁻⁵, 10⁻², the L_2 normalisation factor in [1, 10] and the batch size in [16, 128]. We find the best hyperparameters separately for each prediction problem as reported in § 4.

3.2.2. Random forest

Random forest (Breiman 2001) is a supervised learning algorithm that uses ensemble methods of bootstrap aggregation to solve classification or regression problems. It constructs multiple decision trees during training and then outputs the mean predictions of the individual tree models for testing. It leverages both the bagging method and feature randomness (Breiman 2001) to create an uncorrelated forest of decision trees. At each training step, the model creates a random subset of features and builds simplified trees using these subsets, and while splitting a node, the algorithm searches for the best features among the specific subset. A generalised result then combines the results of numerous uncorrelated models to output the prediction. Because the averaging of uncorrelated tress lowers the overall variance, the random forest model has a lower risk of overfitting and has

been broadly applied to science and engineering domains (Jaiswal & Samikannu 2017; Resende & Drummond 2018).

The hyperparameters of the random forest model include the number of estimators, maximum of tree depth and minimum number of leaves required to split an internal node. To explore the best combinations of settings of the hyperparameters, we use GridSearchCV (Pedregosa *et al.* 2011) with the ranges of [100, 10000], [100, 2000] and [2, 4] for each of the above hyperparameters, respectively. In the following sections, we refer to the random forest regression model as RFR. We present detailed results for this process in § 4.

3.2.3. Gaussian processes regression

The Gaussian processes regression (GPR) model is a probabilistic supervised ML framework that has been widely used for regression and classification tasks. This model can make predictions incorporating prior knowledge and provides uncertainty measures over predictions (Murphy 2013). Similar to other ML models, GPR is classified as a non-parametric model (Liu *et al.* 2020), i.e. when conducting regressions using GPR, the complexity or flexibility of the model is not limited by the number of parameters. More specifically, the model does not assume a specific form for the function that maps the input to the output data, which would be the case for a linear regression model for example.

Formally, a Gaussian process mode is a probability distribution over possible functions that fit a set of points. The regression function modelled by a multivariate Gaussian is given as

$$P(f|X) = \mathcal{N}(f|\mu, K), \tag{3.7}$$

where $X = [x_1, \ldots, x_n]$, $f = [f(x_1), \ldots, f(x_n)]$, $\mu = [m(x_1), \ldots, m(x_n)]$ and $K_{ij} = k(x_i, x_j)$. Here X are the observed data points, m represents the mean function and k represents a positive-definite kernel function. With no observation, the mean function is default to be m(X) = 0 given that the data are often normalised to a zero mean. The Gaussian processes model is a distribution over functions whose shape is defined by K. If points x_i and x_j are considered to be similar by the kernel, function outputs of the two points, $f(x_i)$ and $f(x_j)$, are expected to be similar. In the regression process, given the observed data and a mean function f estimated by these observed points, predictions can be made at new points X_* as $f(X_*)$. Based on the definitions, one can directly derive the joint distribution of f and f_* based on K and, more importantly, the conditional distribution $P(f_*|f,X,X_*)$ over f_* , which is needed for the regression process. These derivations are explained in detail in earlier work (Rasmussen & Williams 2005).

The hyperparameters of the GPR model consist of the kernel and its lengthscale and variance. In this work, we explore six different kernels: Exponential, Matern12, Matern32, Matern52, SquaredExponential and Rational Quadratic. Each of these kernels is added together with a Linear kernel. Generally speaking, the lengthscale controls the smoothness of the learnt functions and the variance controls the scale of values of the learnt functions. Given that the input dimension for our problem is nine, the dimension of the lengthscale is also nine. To optimise these hyperparameters, we perform a grid search using the interval [0.25, 2.1] for each lengthscale value for each input dimension, and the interval [0.1, 2.2] for the variance of the kernels. This search is performed for all six different kernels. We present the final results for this optimisation process in § 4.

3.3. Design and implementation

To implement the three proposed surrogate models we use Python and its supporting libraries. The MLP model is implemented using the Pytorch deep learning framework. We use the Adam solver (Kingma & Ba 2014), an extension of stochastic gradient descent

in the implementation of optimising the weights of the MLP, which is computationally efficient and requires less tuning for hyperparameters.

As for the random forest model, we use scikit-learn (Pedregosa *et al.* 2011), more specifically sklearn.ensemble.RandomForestRegressor, which provides the functions fit and predict for training and prediction, respectively.

Finally, to implement the GPR model, we use an existing Python library called GPflow (Matthews *et al.* 2017), which is based on GPy (Sheffield Machine Learning Software 2012). Although these two libraries are very similar to each other, GPflow leverages TensorFlow for faster computation and gradient calculation, and focuses on variational inference and Markov Chain Monte Carlo methods. One special advantage about this library pertinent to our work is its recent extension to inter-domain approximations and multi-output priors (van der Wilk *et al.* 2020), which is the case of our dataset. As mentioned previously, one essential component of a GPR model is the kernel function. GPflow provides implementations for the most well-known kernels used in the literature, including Exponential, Linear, Matern family, and others. The kernel selection is included in the hyperparameter optimisation process explained in § 3. For the optimisation process, we use the model gpflow.models.GPR also implemented in GPflow, along with gpflow.optimizers.Scipy for function minimisation.

4. Results and evaluation

Based on the designing process and hyperparameter optimisation explained in the previous section, we obtain the best three final ML models to be used as surrogates for the forward modelling problem, predicting current profiles and power deposition. This section presents the details of the final ML models along with their computational performance and prediction accuracy.

4.1. Computational environment and methodology

To create the database of simulations for training and testing the ML surrogate models, we used GENRAY version 10.10_170301 and CQL3D version cswim $_170101.1$ with the Intel FORTRAN compiler version 19.1.3.304 built and run on the MIT Engaging computing cluster at the Massachusetts Green High Performance Computing Center. The MIT PSFC partition of Engaging consists of 136 nodes, with 2×16 Intel Xeon 2.1 GHz cores 128 GB RAM per node. The source code for GENRAY and CQL3D are available at https://github.com/compxco.

Our ML experiments were performed on Cori, a multi-core platform maintained by the National Energy Research Scientific Computing Center (NERSC). All tests were run using Jupyter (https://docs.nersc.gov/services/jupyter/) on a shared CPU node, which consists of a large memory node running an AMD EPYC 7302 3.0 GHz with 32 physical cores, 2 threads per core, 2 sockets per node and has 2 TB of memory (https://docs.nersc.gov/systems/cori/). Appendix A presents more details specifically related to the Python environment used for reproducibility purposes.

The results reported in this section are based on the dataset filtered with the slide-away regime threshold as described in § 3.1. We randomly divide this database of simulations into two subsets: the training set and the test set. The training set is the largest and contains 80% of the overall dataset (10677 samples), whereas the test set contains 20% (2670 samples). The training set is used for both hyperparameter optimisation and to build the final regression model used to predict samples in the test data as described in the next section. For all the models we report execution time for both training and testing, and prediction accuracy based on the MSE.

Problem	H1	H2	НЗ	H4	LR	L2	BATCHSIZE
Curtor	539	254	255	570	3.265	4.051	22
Power	344	601	826	944	2.850	4.873	95

TABLE 2. Best hyperparameters obtained for the MLP model.

4.2. Results of ML model performance and prediction accuracy

4.2.1. Hyperparameters

The first step in the design of the ML models is to optimise their hyperparameters, which is necessary for the selection of the best model for final predictions. In order to find the best hyperparameter set for each model, as described in § 3, the training set is used in a 5-fold cross-validation process using the MSE metric. For each fold, 8541 and 2136 samples are used for train and validation, respectively, and the same 5 folds are used across all 3 models. A separate set of hyperparameters is obtained for each prediction problem (wave power deposition and current profiles).

After the random search approach using the parameters' ranges described in § 3, we find the best set of hyperparameters for the MLP model as presented in table 2. Owing to an existing bigger variation in the power profiles, we observe a denser network is selected with a smaller learning rate and larger batch size in the mini-batch gradient descent for convergence through training.

In figure 2, we report the hyperparameter selection of the random forest model for each prediction problem. The grid search results show that with an increase in the number of estimators, the cross-validated average MSE decreases. The cross-validated accuracy is not sensitive to the max depth of trees trained, for both the current and power profiles. Figure 2(d) presents a minuscule reduction of 5-fold validated accuracy when increasing the minimum sample split in a node. Therefore, we choose 2000 estimators, max depth of 100 and minimum samples split of 2 to balance the computational time and accuracy.

Based on the grid search process using the different hyperparameters (with their respective ranges) related to the GPR model, we obtain a general overview of the cross-validated MSE values across all the 1200 different sets of hyperparameters as shown in figure 3. Based on these results the GPR model presented a low cross-validated MSE variance throughout the hyperparameter optimisation process. For the problem of predicting current profiles, the best model used the Matern32 kernel with 0.75 length scale value for all input dimensions, and variance 0.1. The associated Linear kernel used variance value of 0.6. In the case of predicting power deposition, the best model used the Exponential kernel with 1.25 length scale value for all input dimensions and variance 2.1. The associated Linear kernel used variance value of 1.6. In the following, we present the prediction results using the three ML models obtained based on the best set of hyperparameters.

4.2.2. Prediction results

Once the best hyperparameters are obtained, the models (based on these parameters) are retrained using the whole training data. Then, the final predictions are applied to the test set, which we will refer to as the 'holdout' set, as it includes samples from the simulation database never 'seen' by the models during either hyperparameter optimisation or training. The loss function used in all cases is the MSE, which is used to evaluate and compare our methods.

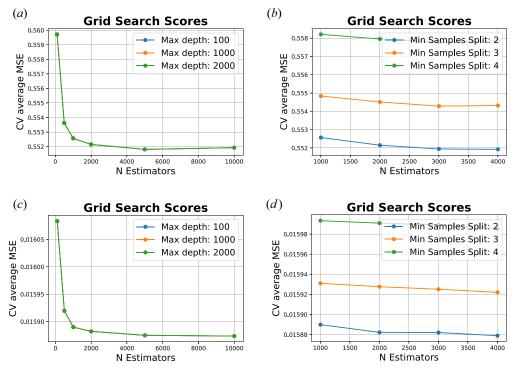


FIGURE 2. Hyperparameter tuning of the random forest regression model for the both current and power profiles: (a,c) cross-validated MSE versus the number of estimators using the maximum tree depth of 100, 1000 and 2000; (b,d) cross-validated MSE versus number of estimators using the minimum samples split of 2, 3 and 4.

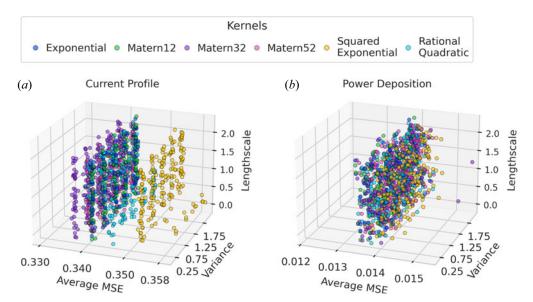


FIGURE 3. Results from the hyperparameter optimisation for the GPR models for both current profile and power deposition. For the current profile the Matern 32 kernel obtained superior results. In the case of power deposition the Exponential kernel performed the best, although the results are very similar among the different set of parameters and kernels.

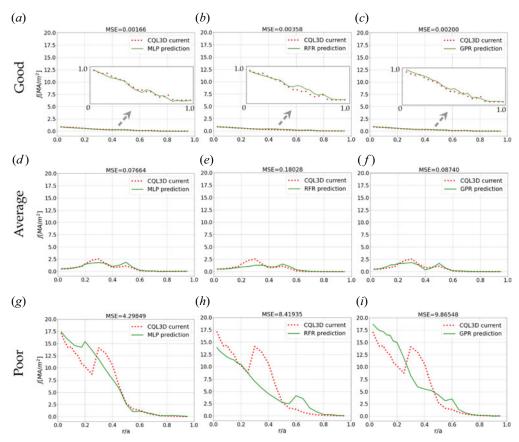


FIGURE 4. Prediction results for current profiles: (a,d,g) MLP, (b,e,h) RFR and (c,f,i) GPR. (a-c) Good fits, (d-f) average fitsand (g-i) poor fits representations for the same test data using the three ML models.

We start by presenting in figure 4 the results for predicting current profiles using the three ML models. Each column represents the ML model used and each row presents typical prediction cases for good fits, average fits, and poor fits. Our methodology for deciding which solution numbers to use for typical good, average, and poor fits is as follows. First, we ordered all prediction results for all methods by MSE. Then, from each set of results ordered by MSE, we considered a relatively small window size of solutions (N = 30) with the lowest MSE and located the solution number common across this set. A similar methodology applies to those with the highest MSE. For the typical average MSE case, due to a high degree of variance in the middle of the MSE distribution, we used a larger window size (N = 160) and display the solution number from the set of mid-range MSE values.

This set of results shows that all three ML methods give an accurate quantitative prediction for the CQL3D current profile in cases with 'good' MSE <0.005. In the case of 'average' MSE of \sim 0.1, the overall magnitude and peak structure of the ML inferences are in excellent agreement with the CQL3D ground truth. The differences between the ML model and the ground truth are within the level of experimental uncertainty in measuring the current profile, and also smaller than typical discrepancy between modelling and

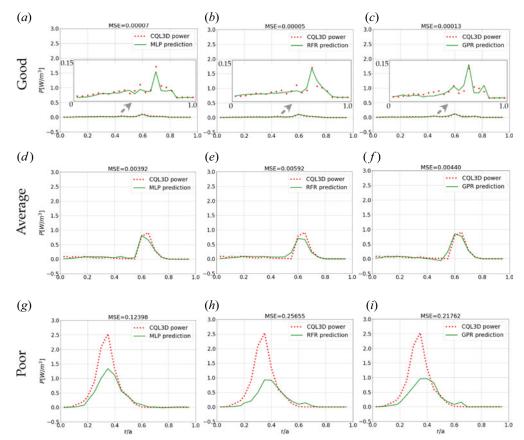


FIGURE 5. Prediction results for power deposition: (a,d,g) MLP, (b,e,h) RFR and (c,f,i) GPR. (a-c) Good fits, (d-f) average fits and (g-i) poor fits representations for the same test data using the three ML models.

experiment. Only in the cases of worst disagreement between the ML model and the ground truth CQL3D model are qualitative differences observable between the two.

Similarly, Figure 5 presents the results for predicting wave power deposition. In the case of the power deposition profile, quantitative agreement is excellent for both the 'good' and 'average' cases, while the 'poor' case infers the correct location for the peak but with a magnitude that is too low by a factor of two. As with the current profiles, the level of agreement between the ML model and the ground truth is well within experimental uncertainty for the average case.

Figures 4 and 5 only show representative prediction examples of 'good', 'average' and 'poor' cases for each ML model. However, in order to show the overall distribution of MSE values for the entire 'holdout' set containing 2670 samples, we present in figure 6 the histograms of MSE values for the three ML models for each prediction problem. Based on these histograms, it is clear that the vast majority of predictions leads to small MSE values, around 10^{-1} for the current profiles and 10^{-2} for the power deposition. Indeed, for the low error regime only MLP has MSE $<10^{-3}$ for the current profiles, and only MLP and RFR have MSE $<5\times10^{-4}$ for the power deposition. MLP has a relative smaller distribution in the high-error regime and outperforms RFR and GPR in both cases. It is interesting to note that even though RFR has a larger distribution of MSE toward the very low-error

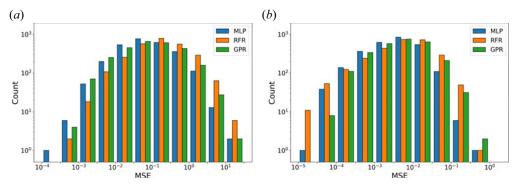


FIGURE 6. Histograms of the MSE in MLP, RFR and GPR ML models for (a) current and (b) power deposition.

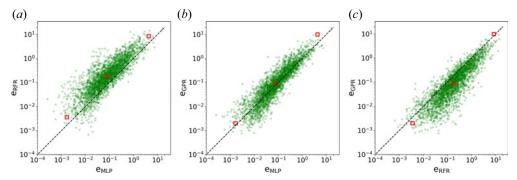


FIGURE 7. Pairwise comparisons of the entire MSE of current prediction among the MLP, RFR and GPR ML models for the inference records. The three representative good, average and poor test instances as in figure 4 are highlighted.

regime compared with the other two models, it also shows a larger MSE at the high-error regimes for the power deposition prediction. For the current profile prediction, GPR in general has more predictions in the lower-error regime indicating a better performance in current profile prediction than RFR.

Figures 7 and 8 compare the MSE of the current profile and power deposition using the three ML models for all the test examples. The highlighted spots correspond to the representative examples in figure 4 and 5. In general, the predictions of the three methods show a strong linear correlation. From a statistical perspective, MLP tends to have a smaller MSE than RFR and GPR in both the low- and high-error regimes for both the prediction cases. GPR shows a very comparable performance as MLP in the low-error regime of the current prediction. For the power prediction, RFR shows a cluster of relatively better performance in the low-error regime, which matches what we observed in the histogram plot of figure 6.

4.3. Discussion

Figures 4 and 5 show anecdotally that the accuracy of all three ML models is anti-correlated with the overall magnitude of the profile. This is potentially problematic as the model is of most interest in the regime where there is significant driven current and power absorption. The origin of this behaviour lies in the loss function used in the

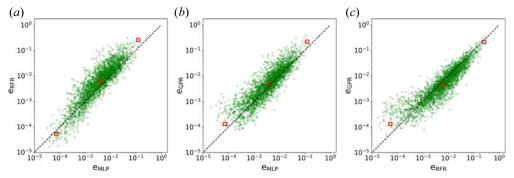


FIGURE 8. Pairwise comparisons of the entire MSE of power prediction among the MLP, RFR and GPR ML models for the inference records. The three representative good, average and poor test instances as in figure 5 are highlighted.

training of the ML models: MSE. The absolute MSE is naturally quite small for signals of small magnitude, whereas for larger-magnitude signals it is naturally larger. This is a double-edged sword. A loss function biased against large-magnitude signals actually improves the accuracy of the ML models in the large-magnitude regime as compared with a normalised loss function, which would weight the large- and small-magnitude signals equally in the training process. Training with a physics-based loss function, which is biased even more strongly against large-magnitude signals, may improve model performance further in the large-magnitude regime.

The previous section shows that the three ML-based surrogate models accurately reproduce the results of GENRAY/CQL3D across the range of inputs in the training database. The ultimate test of the ML-based surrogate models is whether they can produce inferences that are useful. To this end, we compare the MLP-based surrogate model with GENRAY/CQL3D simulations that use a full set of experimental inputs rather than the reduced set of nine zero-dimensional inputs. Figure 9 shows the inferred current density profile for a set of four well-diagnosed EAST discharges. The profiles in this figure use the MLP surrogate model, although the RFR and GPR inferences are substantially similar. The predicted current profiles from the surrogate model reproduce the same qualitative results shown in figure 11(b) of Garofalo et al. (2017), which used full profiles for the plasma density and temperature, a polarimeter-constrained equilibrium reconstruction, and LHCD power from antennas at 2.45 and 4.6 GHz. The ML surrogate model reproduces the current profile peak location accurately, although the magnitude of the current profile peak agrees less closely with the stand-alone GENRAY/CQL3D model in the previously published study. Note that these differences are a result of the reduced number of inputs (e.g. parabolic n_e and T_e profiles as compared with the measured profiles, analytic plasma equilibrium as compared with EFIT reconstruction) to the GENRAY/CQL3D model used to generate the training database, and not inherent to the MLP-based surrogate model. Furthermore, the simulations in figure 9(b) include two frequencies of lower hybrid waves (2.45 and 4.6 GHz), whereas the surrogate model training data is for 4.6 GHz only. Neither figure 9(a) nor figure 9(b) include the effects of fast electron radial diffusion, which are known to be important in reproducing the experimental current profile on EAST (see the differences in figure 11(a-c) in Garofalo *et al.* 2017).

The obtained prediction results clearly show that it is feasible to use ML methods as surrogates for the forward modelling problem achieving high prediction accuracy. Table 3 presents the overall accuracy results for the five-fold cross-validation and testing.

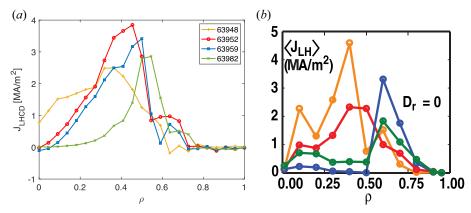


FIGURE 9. (a) Inferred current density profile using the MLP surrogate model for four well-diagnosed EAST discharges. (b) Current profiles for the same set of discharges published in a study by Garofalo. Data in subfigure (b) reproduced from figure 11(b) of Garofalo et al. (2017). The orange-red-blue-green legend applies to both figures.

		5-Fol	Test (hold-out data)			
	Cur	rent	Pov	vers	Current	Powers
Method	$\mu(MSE)$	σ(MSE)	$\mu(MSE)$	σ(MSE)	$\mu(MSE)$	$\mu(MSE)$
Multilayer Perceptron	0.223	0.015	0.009	0.0007	0.227	0.008
Random Forest	0.559	0.045	0.016	0.0007	0.528	0.015
Gaussian Process Regression	0.310	0.019	0.012	0.0006	0.306	0.012

TABLE 3. Evaluation of our three ML models using MSE. For the five-fold cross-validation (CV) process we present the mean (μ) and standard deviation (σ) of the MSE across all folds. The second column presents the prediction results of each final model trained using the full training data.

Based on these results we can conclude that the three models are very stable throughout the cross-validation process, presenting very low standard deviation values for all cases. Moreover, all the three models are accurate in predicting current profiles and power deposition, with MLP and GPR presenting slightly better results compared with random forest.

An additional important point of consideration is the computational cost to achieve such results. Ideally, we would like to use these ML models to perform inference tasks in a fraction of the time required via traditional methods. We report the computational costs required for all the three ML models for both training and inference in table 4. The GPR model requires the highest amount of time to perform the five-fold cross-validation, taking 627.5 and 771.1 min for current and power profiles, respectively. Random forest is the cheapest model, requiring only 3.2 and 1.6 min for each five-fold cross-validation task. Similar computational costs are obtained to train the final models based on the best set of hyperparameters. Most interestingly, all the three models present incredibly low inference times, just under 1 ms for MLP and random forest, and just above 2 ms for GPR, if compared with the 10 min required by the traditional method. This low computational costs easily allow for the applicability of these ML models in real experimental scenarios.

	5-Fold CV		Final model			
	Training time (min)		Training time (min)		Inference time (ms)	
Method	Current	Powers	Current	Powers	Current	Powers
GENRAY/CQL3D	_	_	_	_	10 (min)	
Multilayer perceptron	71.7	119.7	14.4	23.1	0.80	1.69
Random forest	3.2	1.6	9.1	3.9	0.93	0.75
Gaussian process regression	627.5	771.1	508.8	353.9	2.10	2.36

TABLE 4. Timing requirements. The first column reports the total execution time needed to run the entire five-fold cross-validation (CV). The second column reports the time necessary to train each final model along with per sample inference time.

5. Conclusion and future work

The results of this paper demonstrate the feasibility of ML surrogate models for predicting the current profiles and power deposition of GENRAY/CQL3D, especially at the peak locations. The MLP and GPR generate more accurate estimates than RFR in terms of MSE for high-error regimes. On the other hand, RFR is a much more computationally efficient model that requires orders of magnitude less training time, even compared with the cost of one inference in the GENRAY/CQL3D simulation. Still, the computational resources needed to generate the training dataset far exceed the resources needed to train any of the ML surrogate models. Areas with large MSE are also near boundaries of model applicability and where experiments would not typically operate. Future work will include the study and development of customised loss functions based on the physics of the problem to replace MSE as a figure of merit. Similarly, we will work on improving the GPR model to use physics-based kernels aiming for higher prediction accuracy and more efficient training.

A surrogate model is only as accurate as the training data used to create it. The model presented here uses only a small subset of the many hundreds of possible input variables in GENRAY/CQL3D. Increasing the number of input variables will result in more accurate surrogate models at the cost of additional computational resources needed to create a larger training dataset. Even with a small set of zero-dimensional input parameters, the surrogate models reproduce the qualitative results obtained by GENRAY/CQL3D using full experimental profiles.

As noted in § 4, the effects of fast electron radial diffusion are important to reproducing the current profile in EAST with GENRAY/CQL3D. Including radial diffusion in CQL3D increases the computation time by a factor of 50–100. This represents both a challenge and a great opportunity for the ML surrogate models. On the one hand, generating the training database for the surrogate model is significantly more expensive when radial diffusion is included, whereas the time savings as compared with the original full-fidelity model are enhanced by an even greater factor. We hypothesise that it may be possible to reduce the size of the training database somewhat because radial diffusion smooths out the current profiles, which may allow for low MSE with a smaller database of simulations.

Ultimately, ML-based fast surrogate models will be incorporated into larger frameworks for plasma prediction and control. Control system development frameworks such as COTSIM (Pajares & Schuster 2019, 2020) and integrated models such as TRANSP (Jardin, Pomphrey & Delucia 1986; Wehner *et al.* 2019) or IPS-FASTRAN (Elwasif *et al.* 2010;

Park *et al.* 2017) will benefit from more rapid estimates of power deposition and current drive profiles. To facilitate adoption of these surrogate models, we have made the trained models available online through Dataverse, as described in Appendix A. Rapid iteration times will allow increased exploration of parameter space to find optimal operating points using plasma 'flight simulators' prior to experiments, or in the design of new fusion reactors.

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Declaration of interests

The authors report no conflict of interest.

Data availability statement

The data that support the findings of this study are openly available in Dataverse at https://doi.org/10.7910/DVN/5YY6PE and in GitHub at https://github.com/FusionRFModelingMachineLearning/genray-cql3d-ml. See additional details in Appendix A.

Appendix A. Reproducibility

This appendix provides information about accessing and using source code, trained models and sample data that will facilitate reproducing the work we have presented earlier in the paper. Our team is providing four sets of artifacts to facilitate reproducibility and potential application to additional fusion devices. These consist of the final trained models for each of the regression methods used (appendix A.1), the data for performing both training and inference using these methods (appendix A.2), source code to train the ML models using the same dataset used in this article (appendix A.3) and the source code for loading the trained models and sample data and then performing the inference computations (appendix A.4), as described in the following.

A.1. Trained models in ONNX and tensorflow formats

We leverage ONNX as much as possible for sharing our trained ML models. ONNX is a storage format and set of utilities for storing and loading ML models to/from disk files (Bai *et al.* 2019). The ONNX project is a collaboration of numerous partners from industry.

For each the MLP and RFR models we describe earlier (§ 3), we have saved each of these trained models in ONNX format. Owing to current limitations on the operations supported by ONNX, we are making the GPR trained model available using the Tensorflow format. Each trained model reflects the hyperparameter optimisation and training process we

describe earlier (§ 3). The inference results we obtain (§ 4) come from these three trained models prior to export to ONNX and Tensorflow formats.¹

A.2. Data comprising database of simulations, training dataset and sample data for inference

We are providing several sets of data via Dataverse (https://doi.org/10.7910/DVN/5YY6PE). First is the set of NETCDF outputs generated for the full database of simulations, including additional code outputs (e.g. hard X-ray emission) not used in this work. Second is the dataset filtered with the slide-away regime threshold consisting of 13 347 samples, which can be used to train the three ML models described in this article using CSV format. The third is a sample dataset that can be leveraged to perform inference using the provided trained models in ONNX and Tensorflow using CSV format.

A.3. Python scripts for training ML models

The source code used to train the three ML models presented here are available in our GitHub repository (https://github.com/FusionRFModelingMachineLearning/genray-cql3d-ml). These scripts allow for retraining the final models with the dataset containing 13 347 samples using the best parameters found via hyperparameter optimisation. Detailed information about the required computational environment and instructions to perform this tasks are available in the repository.

A.4. Python scripts for loading models, data and performing inference workload

In addition to the source code to train the ML models, we also provide Python scripts that directly load the trained models in ONNX and Tensorflow formats in order to perform inference with the provided sample data. We provide separate scripts for each one of the models as described in details in our GitHub repository.

REFERENCES

- ABBATE, J., CONLIN, R. & KOLEMEN, E. 2021 Data-driven profile prediction for DIII-D. *Nucl. Fusion* **61** (4), 046027.
- ALLEN, L. & BISHOP, C.M. 1992 Neural network approach to energy confinement scaling in tokamaks. *Plasma Phys. Control. Fusion* **34** (7), 1291.
- BAI, J., LU, F. & ZHANG, K. 2019 Onnx: open neural network exchange. https://github.com/onnx/onnx.
- BAI, Z. & PENG, L. 2021 Non-intrusive nonlinear model reduction via machine learning approximations to low-dimensional operators. *Adv. Model. Simul. Engng Sci.* 8 (1), 1–24.
- BISHOP, C., COX, P., HAYNES, P., ROACH, C., SMITH, M., TODD, T. & TROTMAN, D. 1992 A neural network approach to tokamak equilibrium control. In *Neural Network Applications* (ed. Taylor J.G.), pp. 114–128. Springer.
- BISHOP, C.M. 1995b Neural Networks for Pattern Recognition. Oxford University Press.
- BISHOP, C.M., HAYNES, P.S., SMITH, M.E.U., TODD, T.N. & TROTMAN, D.L. 1995a Real-time control of a tokamak plasma using neural networks. *Neural Comput.* 7 (1), 206–217.
- BISHOP, C., STRACHAN, I., O'ROURKE, J., MADDISON, G. & THOMAS, P. 1993 Reconstruction of tokamak density profiles using feedforward networks. *Neural Comput. Appl.* 1 (1), 4–16.
- BOYER, M.D., KAYE, S. & ERICKSON, K. 2019 Real-time capable modeling of neutral beam injection on NSTX-u using neural networks. *Nucl. Fusion* **59** (5), 056008.
- BRAMBILLA, M. 1996 A full wave code for ion cyclotron waves in toroidal plasmas, IPP Report 5/66. Max-Planck-Institut für Plasmaphysik.
- BREIMAN, L. 2001 Random forests. Mach. Learn. 45 (1), 5-32.

¹Tensorflow file format, online at https://www.tensorflow.org/hub/model_formats, last accessed May 2022.

- BRUNTON, S.L. & KUTZ, J.N. 2022 Data-Driven Science and Engineering: Machine Learning, Dynamical Systems, and Control. Cambridge University Press.
- CHURCHILL, R.M., CHOI, J., KUBE, R., CHANG, C.S. & KLASKY, S. 2020b Machine learning for the complex, multi-scale datasets in fusion energy. In *Smoky Mountains Computational Sciences and Engineering Conference* (eds. Nichols J., Verastegui B., Maccabe A., Hernandez O., Parete-Koon S. & Ahearn T.), pp. 269–284. Springer.
- CHURCHILL, R.M., TOBIAS, B., ZHU, Y. & DIII-D TEAM 2020a Deep convolutional neural networks for multi-scale time-series classification and application to tokamak disruption prediction using raw, high temporal resolution diagnostic data. *Phys. Plasmas* 27 (6), 062510.
- CONNOR, J.W. & HASTIE, R.J. 1975 Relativistic limitations on runaway electrons. *Nucl. Fusion* **15** (3), 415.
- DREICER, H. 1959 Electron and ion runaway in a fully ionized gas. I. Phys. Rev. 115 (2), 238.
- DURAISAMY, K., IACCARINO, G. & XIAO, H. 2019 Turbulence modeling in the age of data. *Annu. Rev. Fluid Mech.* **51**, 357–377.
- ELWASIF, W.R., BERNHOLDT, D.E., SHET, A.G., FOLEY, S.S., BRAMLEY, R., BATCHELOR, D.B. & BERRY, L.A. 2010 The design and implementation of the swim integrated plasma simulator. In 2010 18th Euromicro Conference on Parallel, Distributed and Network-based Processing, pp. 419–427. IEEE.
- GALAMBOS, J.D., PERKINS, L.J., HANEY, S.W. & MANDREKAS, J. 1995 Commercial tokamak reactor potential with advanced tokamak operation. *Nucl. Fusion* **35** (5), 551.
- GARDILL, M., FUCHS, J., FRANK, C. & WEIGEL, R. 2018 A multi-layer perceptron applied to number of target indication for direction-of-arrival estimation in automotive radar sensors. In 2018 IEEE 28th International Workshop on Machine Learning for Signal Processing (MLSP), pp. 1–6. IEEE.
- GAROFALO, A.M., GONG, X.Z., QIAN, J., CHEN, J., LI, G., LI, K., LI, M.H., ZHAI, X., BONOLI, P., BROWER, D., *et al.* 2017 Development of high poloidal beta, steady-state scenario with ITER-like tungsten divertor on EAST. *Nucl. Fusion* 57 (7), 076037.
- GOLDSTON, R.J. 1996 Physics of the steady-state advanced tokamak. Phys. Plasmas 3 (5), 1794–1802.
- GUAZZOTTO, L. & FREIDBERG, J.P. 2021 Simple, general, realistic, robust, analytic tokamak equilibria. Part 1. Limiter and divertor tokamaks. *J. Plasma Phys.* 87 (3), 905870303.
- HARVEY, R.W. & McCoy, M. 1992 The CQL3D Fokker–Planck Code. In *Proceedings of the IAEA Technical Committee Meeting on Simulation and Modeling of Thermonuclear Plasmas*, pp. 489–526.
- HATFIELD, P.W., GAFFNEY, J.A., ANDERSON, G.J., ALI, S., ANTONELLI, L., BAŞEĞMEZ DU PREE, S., CITRIN, J., FAJARDO, M., KNAPP, P., KETTLE, B., *et al.* 2021 The data-driven future of high-energy-density physics. *Nature* **593** (7859), 351–361.
- HAYASHI, N. 2010 Advanced tokamak research with integrated modeling in JT-60 upgrade. *Phys. Plasmas* 17 (5), 056112.
- HERNANDEZ, J.V., VANNUCCI, A., TAJIMA, T., LIN, Z., HORTON, W. & MCCOOL, S.C. 1996 Neural network prediction of some classes of tokamak disruptions. *Nucl. Fusion* **36** (8), 1009.
- HOULBERG, W.A. & ATTENBERGER, S.E. 1994 Evaluation of current drive requirements and operating characteristics of a high bootstrap fraction advanced tokamak reactor. *Fusion Technol.* **26** (3P2), 566–571.
- JAISWAL, J.K. & SAMIKANNU, R. 2017 Application of random forest algorithm on feature subset selection and classification and regression. In 2017 World Congress on Computing and Communication Technologies (WCCCT), pp. 65–68. IEEE.
- JARDIN, S.C., POMPHREY, N. & DELUCIA, J. 1986 Dynamic modeling of transport and positional control of tokamaks. *J. Comput. Phys.* **66** (2), 481–507.
- KIKUCHI, M. 1993 Prospects of a stationary tokamak reactor. *Plasma Phys. Control. Fusion* **35** (SB), B39. KINGMA, D.P. & BA, J. 2014 Adam: a method for stochastic optimization. arXiv:1412.6980.
- LISTER, J.B. & SCHNURRENBERGER, H. 1991 Fast non-linear extraction of plasma equilibrium parameters using a neural network mapping. *Nucl. Fusion* **31** (7), 1291.

- LIU, F.K., DING, B.J., LI, J.G., WAN, B.N., SHAN, J.F., WANG, M., LIU, L., ZHAO, L.M., LI, M.H., LI, Y.C., *et al.* 2015 First results of LHCD experiments with 4.6 GHz system toward steady-state plasma in east. *Nucl. Fusion* 55 (12), 123022.
- LIU, H., ONG, Y., SHEN, X. & CAI, J. 2020 When gaussian process meets big data: a review of scalable GPs. *IEEE Trans. Neural Netw. Learn. Syst.* 31, 4405–4423.
- LUCE, T.C. 2011 Realizing steady-state tokamak operation for fusion energy. *Phys. Plasmas* 18 (3), 030501.
- MALKIEL, I., MREJEN, M., NAGLER, A., ARIELI, U., WOLF, L. & SUCHOWSKI, H. 2018 Plasmonic nanostructure design and characterization via deep learning. *Light: Sci. Appl.* 7, 60.
- MASCHLER, B. & WEYRICH, M. 2021 Deep transfer learning for industrial automation: a review and discussion of new techniques for data-driven machine learning. *IEEE Ind. Electron. Mag.* **15** (2), 65–75.
- MATTHEWS, A.G.G., VAN DER WILK, M., NICKSON, T., FUJII, K., BOUKOUVALAS, A., LEÓN-VILLAGRÁ, P., GHAHRAMANI, Z. & HENSMAN, J. 2017 GPflow: a Gaussian process library using TensorFlow. *J. Machine Learning Res.* 18 (40), 1–6.
- MCKAY, M.D., CONOVER, W.J. & WHITEMAN, D.E. 1976 Report on the application of statistical techniques to the analysis of computer codes. *Tech. Rep.* Los Alamos Scientific Lab., N. Mex.(USA).
- MENEGHINI, O., SMITH, S.P., LAO, L.L., IZACARD, O., REN, Q., PARK, J.M., CANDY, J., WANG, Z., LUNA, C.J., IZZO, V.A., *et al.* 2015 Integrated modeling applications for tokamak experiments with OMFIT. *Nucl. Fusion* 55 (8), 083008.
- MENEGHINI, O., SNOEP, G., LYONS, B.C., MCCLENAGHAN, J., IMAI, C.S., GRIERSON, B., SMITH, S.P., STAEBLER, G.M., SNYDER, P.B., CANDY, J., *et al.* 2020 Neural-network accelerated coupled core-pedestal simulations with self-consistent transport of impurities and compatible with ITER IMAS. *Nucl. Fusion* **61** (2), 026006.
- MILLER, M.A., CHURCHILL, R.M., DENER, A., CHANG, C.S., MUNSON, T. & HAGER, R. 2021 Encoder–decoder neural network for solving the nonlinear Fokker–Planck–Landau collision operator in XGC. *J. Plasma Phys.* 87 (2), 905870211.
- MORABITO, F.C. & VERSACI, M. 1997 A fuzzy-neural approach to real time plasma boundary reconstruction in tokamak reactors. In *Proceedings of International Conference on Neural Networks* (ICNN'97), vol. 1, pp. 43–47. IEEE.
- MOREAU, D. & VOITSEKHOVITCH, I. 1999 Plasma control issues for an advanced steady state tokamak reactor. *Nucl. Fusion* **39** (5), 685.
- MOREAU, D., WALKER, M.L., FERRON, J.R., LIU, F., SCHUSTER, E., BARTON, J.E., BOYER, M.D., BURRELL, K.H., FLANAGAN, S.M., GOHIL, P., et al. 2013 Integrated magnetic and kinetic control of advanced tokamak plasmas on DIII-D based on data-driven models. *Nucl. Fusion* 53 (6), 063020.
- MOROSOHK, S.M., BOYER, M.D. & SCHUSTER, E. 2021 Accelerated version of NUBEAM capabilities in DIII-D using neural networks. *Fusion Engng Des.* 163, 112125.
- MURPHY, K.P. 2013 Machine Learning: A Probabilistic Perspective. MIT Press.
- NAJMABADI, F., ABDOU, A., BROMBERG, L., BROWN, T., CHAN, V.C., CHU, M.C., DAHLGREN, F., EL-GUEBALY, L., HEITZENROEDER, P., HENDERSON, D., *et al.* 2006 The ARIES-AT advanced tokamak, advanced technology fusion power plant. *Fusion Engng Des.* **80** (1), 3–23. ARIES AT Special Issue.
- PAJARES, A. & SCHUSTER, E. 2019 Actuator management via real-time optimization for integrated control in tokamaks. In *Proc. 46th EPS Conf. Plasma Phys.*, p. 1.
- PAJARES, A. & SCHUSTER, E. 2020 Integrated robust control of the global toroidal rotation and total plasma energy in tokamaks. *IEEE Trans. Plasma Sci.* **48** (6), 1606–1612.
- PARK, J.M., MURAKAMI, M., ST. JOHN, H.E., LAO, L.L., CHU, M.S. & PRATER, R. 2017 An efficient transport solver for tokamak plasmas. *Comput. Phys. Commun.* 214, 1–5.
- PEDREGOSA, F., VAROQUAUX, G., GRAMFORT, A., MICHEL, V., THIRION, B., GRISEL, O., BLONDEL, M., PRETTENHOFER, P., WEISS, R., DUBOURG, V., et al. 2011 Scikit-learn: machine learning in Python. *J. Machine Learning Res.* 12, 2825–2830.
- PETROV, Y.V. & HARVEY, R.W. 2016 A fully-neoclassical finite-orbit-width version of the CQL3D Fokker–Planck code. *Plasma Phys. Control. Fusion* **58** (11), 115001.

- POLI, F.M. 2018 Integrated tokamak modeling: when physics informs engineering and research planning. *Phys. Plasmas* **25** (5), 055602.
- RASMUSSEN, C.E. & WILLIAMS, C.K.I. 2005 Gaussian Processes for Machine Learning (Adaptive Computation and Machine Learning). The MIT Press.
- REA, C. & GRANETZ, R.S. 2018 Exploratory machine learning studies for disruption prediction using large databases on DIII-D. *Fusion Sci. Technol.* **74** (1–2), 89–100.
- REA, C., MONTES, K.J., PAU, A., GRANETZ, R.S. & SAUTER, O. 2020 Progress toward interpretable machine learning-based disruption predictors across tokamaks. *Fusion Sci. Technol.* **76** (8), 912–924.
- REICHSTEIN, M., CAMPS-VALLS, G., STEVENS, B., JUNG, M., DENZLER, J., CARVALHAIS N. & PRABHAT 2019 Deep learning and process understanding for data-driven Earth system science. *Nature* **566** (7743), 195–204.
- RESENDE, P.A.A. & DRUMMOND, A.C. 2018 A survey of random forest based methods for intrusion detection systems. *ACM Comput. Surv. (CSUR)* **51** (3), 1–36.
- RODRIGUEZ-FERNANDEZ, P., WHITE, A.E., CREELY, A.J., GREENWALD, M.J., HOWARD, N.T., SCIORTINO, F. & WRIGHT, J.C. 2018 VITALS: a surrogate-based optimization framework for the accelerated validation of plasma transport codes. *Fusion Sci. Technol.* **74** (1–2), 65–76.
- RUMELHART, D.E., HINTON, G.E. & WILLIAMS, R.J. 1986 Learning representations by back-propagating errors. *Nature* **323** (6088), 533–536.
- SHEFFIELD MACHINE LEARNING SOFTWARE 2012 GPy: a Gaussian process framework in Python. http://github.com/SheffieldML/GPy.
- SHIRAIWA, S., GREENWALD, M., STILLERMAN, J.A., WALLACE, G.M., LONDON, M.R. & THOMAS, J. 2018 Tools to export published datasets together with metadata from IDL/Python/MATLAB and πScope. *Fusion Engng Des.* **130**, 104–106.
- SHIRAIWA, S., WRIGHT, J.C., LEE, J.P. & BONOLI, P.T. 2017 HIS-TORIC: extending core ICRF wave simulation to include realistic SOL plasmas. *Nucl. Fusion* 57 (8), 086048.
- SIPS, A.C.C. 2005 Advanced scenarios for ITER operation. *Plasma Phys. Control. Fusion* **47** (5A), A19. SMIRNOV, A.P. & HARVEY, R.W. 1995 A general ray tracing code GENRAY. *Bull. Am. Phys. Soc.* **40**, 1837.
- VAN DE PLASSCHE, K.L., CITRIN, J., BOURDELLE, C., CAMENEN, Y., CASSON, F.J., DAGNELIE, V.I., FELICI, F., HO, A., VAN MULDERS, S. & JET CONTRIBUTORS 2020 Fast modeling of turbulent transport in fusion plasmas using neural networks. *Phys. Plasmas* 27 (2), 022310.
- VAN DER WILK, M., DUTORDOIR, V., JOHN, S.T., ARTEMEV, A., ADAM, V. & HENSMAN, J. 2020 A framework for interdomain and multioutput Gaussian processes. arXiv:2003.01115.
- WALLACE, G.M., PARKER, R.R., BONOLI, P.T., HUBBARD, A.E., HUGHES, J.W., LABOMBARD, B.L., MENEGHINI, O., SCHMIDT, A.E., SHIRAIWA, S., WHYTE, D.G., *et al.* 2010 Absorption of lower hybrid waves in the scrape off layer of a diverted tokamak. *Phys. Plasmas* 17 (8), 082508.
- WEHNER, W.P., SCHUSTER, E., BOYER, M.D. & POLI, F.M. 2019 TRANSP-based optimization towards tokamak scenario development. *Fusion Engng Des.* **146**, 547–550, sI:SOFT-30.
- WILKES, E.H., EMMETT, E., BELTRAN, L., WOODWARD, G.M. & CARLING, R.S. 2020 A machine learning approach for the automated interpretation of plasma amino acid profiles. *Clin. Chem.* 66 (9), 1210–1218.
- WRÓBLEWSKI, D. 1997 Neural network evaluation of tokamak current profiles for real time control. *Rev. Sci. Instrum.* **68** (2), 1281–1285.
- WROBLEWSKI, D., JAHNS, G.L. & LEUER, J.A. 1997 Tokamak disruption alarm based on a neural network model of the high-beta limit. *Nucl. Fusion* **37** (6), 725.
- WU, S. & EAST TEAM. 2007 An overview of the east project. Fusion Engng Des. 82 (5-14), 463-471.
- YELLAPANTULA, S., DE FRAHAN, M.T.H., KING, R., DAY, M. & GROUT, R. 2020 Machine learning of combustion les models from reacting direct numerical simulation. In *Data Analysis for Direct Numerical Simulations of Turbulent Combustion* (eds. Pitsch H. & Attili A.), pp. 273–292. Springer.
- YÜKSEL, E., SOYDANER, D. & BAHTIYAR, H. 2021 Nuclear binding energy predictions using neural networks: application of the multilayer perceptron. arXiv:2101.12117.