

Scattering Length and the Spectrum of $-\Delta + V$

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Abstract. Given a non-negative, locally integrable function V on \mathbb{R}^n , we give a necessary and sufficient condition that $-\Delta + V$ have purely discrete spectrum, in terms of the scattering length of V restricted to boxes.

1 Introduction

It is a classical result of K. Friedrichs [F] that if $V \in L^1_{\text{loc}}(\mathbb{R}^n)$ and $V \geq 0$, then $-\Delta + V$ yields a positive self-adjoint operator on $L^2(\mathbb{R}^n)$, and its spectrum is discrete if $V(x) \rightarrow +\infty$ as $|x| \rightarrow \infty$. A. Molchanov [Mol] produced a necessary and sufficient condition for such an operator to have discrete spectrum. His condition takes the form

$$(1.1) \quad \inf_F \int_{Q_{b,\xi} \setminus F} V(x) dx \rightarrow \infty, \quad \text{as } |\xi| \rightarrow \infty,$$

for each $b \in (0, 1]$, where $Q_{b,\xi}$ is the n -dimensional cube of the form

$$(1.2) \quad Q_{b,\xi} = \left\{ x \in \mathbb{R}^n : \xi_j - \frac{b}{2} \leq x_j \leq \xi_j + \frac{b}{2} \right\}.$$

(We henceforth say $Q_{b,\xi}$ is the cube with sidelength b and center ξ .) In (1.1), F runs over the “negligible” subsets of $Q_{b,\xi}$, defined by the condition $\text{cap } F \leq \gamma \text{cap } Q_{b,\xi}$. In [Mol], γ was taken to be a particular (small) constant γ_n .

Recent important work of V. Maz'ya and M. Shubin [MS] provides a cleaner form for the necessary and sufficient condition. In particular, γ can be given any value in $(0, 1)$. Furthermore, they allow $\gamma = \gamma(b)$, possibly decaying to 0 as $b \rightarrow 0$, as long as $b^{-2}\gamma(b) \rightarrow \infty$.

Our purpose here is to produce an alternative formulation of a necessary and sufficient condition that $-\Delta + V$ have discrete spectrum (given $V \geq 0$, $V \in L^1_{\text{loc}}(\mathbb{R}^n)$). Our result is phrased in terms of “scattering length,” a quantity $\Gamma(v)$ associated to integrable $v \geq 0$ that is somewhat parallel to the notion of capacity of a set. In fact, if K is a compact set satisfying a mild regularity condition,

$$(1.3) \quad \text{cap } K = \lim_{r \rightarrow +\infty} \Gamma(r\chi_K),$$

where χ_K denotes the characteristic function of K . We will recall the definition of $\Gamma(v)$ in §2. Our main result is the following.

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Theorem 1.1 Given $V \geq 0$, $V \in L^1_{\text{loc}}(\mathbb{R}^n)$, the following three conditions are equivalent.

- (1) $-\Delta + V$ has purely discrete spectrum on $L^2(\mathbb{R}^n)$.
- (2) Given $A \in (0, \infty)$, there exists $b = b(A) \in (0, 1]$ and $R \in (0, \infty)$ such that

$$\Gamma(b^2V_{b,\xi}) \geq Ab^2, \quad \text{for } |\xi| \geq R.$$

- (3) Given $A \in (0, \infty)$, there exists $b_0 = b_0(A) \in (0, 1]$ and $R: (0, b_0] \rightarrow (0, \infty)$ such that

$$\Gamma(b^2V_{b,\xi}) \geq Ab^2, \quad \text{for } b \in (0, b_0], |\xi| \geq R(b).$$

Here $V_{b,\xi}$ is a positive function supported on the unit cube $Q = Q_{1,0}$, given by

$$(1.4) \quad V_{b,\xi}(x) = V(bx + \xi), \quad x \in Q.$$

The rest of this paper is structured as follows. In §2 we define $\Gamma(\nu)$ for positive, integrable ν and review some of its crucial properties. In §3 we prove that (2) \Rightarrow (1) in Theorem 1.1, and in §4 we prove that (1) \Rightarrow (3). Clearly (3) \Rightarrow (2), so this will prove Theorem 1.1. There is one result in §4, Lemma 4.2, whose proof is presented in §5.

Remark In the formal limit $V = +\infty$ on $K = \mathbb{R}^n \setminus \Omega$, where one considers $-\Delta$ on $L^2(\Omega)$, with the Dirichlet boundary condition on $\partial\Omega$, the condition (3) of Theorem 1.1 becomes that for each $A \in (0, \infty)$, there exists $b_0 = b_0(A) \in (0, 1]$ and $R: (0, b_0] \rightarrow (0, \infty)$ such that

$$(1.5) \quad \text{cap } K_{b,\xi} \geq Ab^2(\text{cap } Q_{b,\xi}), \quad \forall b \in (0, b_0], |\xi| \geq R(b),$$

where $K_{b,\xi} = K \cap Q_{b,\xi}$. This coincides with one of the criteria (necessary and sufficient) for discreteness presented in [MS, Remark 2.7].

2 Scattering Length

Here we define the scattering length $\Gamma(\nu)$ of a positive integrable potential ν and review some of its properties. Our material is taken from [T], which in turn was influenced by results on scattering length presented in [K, KL]. For simplicity we take $n \geq 3$.

To such ν we associate the capacity potential U_ν and the scattering length $\Gamma(\nu)$ as follows. First assume that $\nu \in L^2(\mathbb{R}^n)$ and has support in a compact set K , as well as $\nu \geq 0$. We define U_ν by

$$(2.1) \quad U_\nu(x) = \lim_{\varepsilon \searrow 0} (\varepsilon + \nu - \Delta)^{-1} \nu(x).$$

It is shown that this limit exists in $L^2_{\text{loc}}(\mathbb{R}^n)$ and satisfies

$$(2.2) \quad 0 \leq U_\nu \leq 1, \quad \nu \leq w \Rightarrow U_\nu \leq U_w.$$

The existence proof in [K, KL] involves producing the formula

$$(2.3) \quad U_\nu(x) = E_x \left\{ 1 - \exp \left(- \int_0^\infty \nu(b(\tau)) d\tau \right) \right\},$$

where E_x is expectation with respect to Wiener measure on Brownian paths b starting at x ; see also [T, p. 292] for a derivation of this formula.

The function U_ν solves the PDE

$$(2.4) \quad \Delta U_\nu = -\nu(1 - U_\nu).$$

It follows that $-\Delta U_\nu = \mu_\nu$ is a positive measure on \mathbb{R}^n . We set

$$(2.5) \quad \Gamma(\nu) = \int d\mu_\nu(x).$$

Some basic results on $\Gamma(\nu)$ include:

$$(2.6) \quad \begin{aligned} \nu \leq w &\implies \Gamma(\nu) \leq \Gamma(w), \\ \Gamma(\nu + w) &\leq \Gamma(\nu) + \Gamma(w), \\ \nu_n \nearrow \nu &\implies \Gamma(\nu_n) \nearrow \Gamma(\nu), \\ \Gamma(\nu) &\leq \|\nu\|_{L^1}. \end{aligned}$$

We also have

$$(2.7) \quad \|\nabla U_\nu\|_{L^2(\mathbb{R}^n)}^2 = \int_{\mathbb{R}^n} U_\nu(x) d\mu_\nu(x) \leq \Gamma(\nu),$$

and, for any ball $B \subset \mathbb{R}^n$,

$$(2.8) \quad \|U_\nu\|_{L^1(B)} \leq \alpha(B)\Gamma(\nu).$$

These results are established in [T, Propositions 1.2–1.6]. They allow us to define U_ν and $\Gamma(\nu)$ for positive $\nu \in L^1(\mathbb{R}^n)$, having

$$(2.9) \quad \nu_n \nearrow \nu, \nu_n \in L^2_{\text{comp}}(\mathbb{R}^n) \implies U_{\nu_n} \nearrow U_\nu, \Gamma(\nu_n) \nearrow \Gamma(\nu).$$

We now give two key estimates, established in [T], which connect scattering length to eigenvalue estimates. Suppose $\nu \geq 0$ is an integrable function supported on Q , the cube of sidelength 1 centered at 0. Let $\lambda_1(\nu) \in [0, \infty)$ denote the smallest eigenvalue of $-\Delta + \nu$, with the Neumann boundary condition, on $L^2(Q)$. The following result summarizes [T, Propositions 2.2–2.3].

Proposition 2.1 *There exists $C_n \in (0, \infty)$ such that*

$$(2.10) \quad \lambda_1(\nu) \geq C_n \Gamma(\nu).$$

Furthermore, there exist $E_n, \tilde{C}_n \in (0, \infty)$ such that

$$(2.11) \quad \Gamma(\nu) \leq E_n \implies \lambda_1(\nu) \leq \tilde{C}_n \Gamma(\nu).$$

We refer to [T, pp. 295–297] for proofs of these results. We mention that (2.11) is proven by an apt choice of test function in the variational characterization of $\lambda_1(\nu)$, while (2.10) is proven by examining the decay rate for e^{-tL_N} , where L_N denotes $-\Delta + \nu$, with the Neumann boundary condition, on $L^2(Q)$.

3 Sufficient Condition for Discrete Spectrum

The following result yields the implication (2) \Rightarrow (1) in Theorem 1.1.

Proposition 3.1 *Take $A \in (0, \infty)$ and let C_n be as in (2.10). Assume that there exists $b = b(A) \in (0, 1]$ and $R = R(A) \in (0, \infty)$ such that*

$$(3.1) \quad C_n \Gamma(b^2 V_{b,\xi}) \geq Ab^2, \quad \text{for } |\xi| \geq R.$$

Then

$$(3.2) \quad \text{ess spec}(-\Delta + V) \subset [A, \infty).$$

Proof Let $Q_{b,\xi}$ denote the cube of edge b , center ξ , as in (1.2), and let $L_{b,\xi}$ denote the operator $-\Delta + V$ on $L^2(Q_{b,\xi})$, with the Neumann boundary condition. A standard argument involving Rellich's theorem shows that, if there exists $R = R(A)$ such that

$$(3.3) \quad \text{spec } L_{b,\xi} \subset [A, \infty), \quad \text{for } |\xi| \geq R,$$

then (3.2) holds. Now $L_{b,\xi}$ is unitarily equivalent to the operator

$$(3.4) \quad -b^{-2}\Delta + V_{b,\xi} = b^{-2}(-\Delta + b^2 V_{b,\xi}),$$

on $L^2(Q)$ (Q denoting the cube of edge 1, center 0), where

$$(3.5) \quad V_{b,\xi}(x) = V(bx + \xi), \quad x \in Q,$$

and one places the Neumann boundary condition on the operator (3.4). Now, by Proposition 2.1, the spectrum of this operator is bounded below by

$$(3.6) \quad C_n b^{-2} \Gamma(b^2 V_{b,\xi}),$$

so Proposition 3.1 is proven. ■

4 Necessary Condition for Discrete Spectrum

It is convenient to set up some notation. Given a cube $Q_\nu \subset \mathbb{R}^n$, we denote by

$$(4.1) \quad \lambda_D^{Q_\nu}(-\Delta + V), \quad \text{resp.}, \quad \lambda_N^{Q_\nu}(-\Delta + V),$$

the smallest eigenvalue of $-\Delta + V$ on $L^2(Q_\nu)$, where we impose, respectively, the Dirichlet or Neumann boundary condition on ∂Q_ν . As before, let $Q_{b,\xi}$ denote the cube of edge b , center ξ , as in (1.2). We continue to assume $V \geq 0$ and $V \in L^1_{\text{loc}}(\mathbb{R}^n)$.

Lemma 4.1 *If $-\Delta + V$ has discrete spectrum on $L^2(\mathbb{R}^n)$, then for each $b \in (0, 1]$,*

$$(4.2) \quad \lambda_D^{Q_{b,\xi}}(-\Delta + V) \longrightarrow +\infty, \quad \text{as } |\xi| \rightarrow \infty.$$

Proof As is well known, $-\Delta + V$ has discrete spectrum on $L^2(\mathbb{R}^n)$ if and only if the set

$$(4.3) \quad X = \{u \in H^1(\mathbb{R}^n) : \|\nabla u\|_{L^2}^2 + \|V^{1/2}u\|_{L^2}^2 \leq 1\}$$

is compact in $L^2(\mathbb{R}^n)$. In turn, such compactness implies

$$(4.4) \quad \int_{|x| \geq R} |u(x)|^2 dx \leq \varepsilon(R), \quad \forall u \in X,$$

where $\varepsilon(R) \rightarrow 0$ as $R \rightarrow \infty$. If we restrict attention to $u \in H_0^1(Q_{b,\xi})$, this gives (4.2). ■

In the following lemma, $Q = Q_{1,0}$, the unit cube centered at 0.

Lemma 4.2 *There exists $A_n \in (0, \infty)$ and $B_n: [A_n, \infty) \rightarrow (0, \infty)$, such that $B_n(A) \rightarrow \infty$ as $A \rightarrow \infty$, and such that whenever $v \in L^1(Q)$ is non-negative and whenever $A \geq A_n$,*

$$(4.5) \quad \lambda_D^Q(-\Delta + v) \geq A \implies \lambda_N^Q(-\Delta + v) \geq B_n(A).$$

Such a result is established in [Mol]; a proof is also given in [KS, Lemma 2.9]. For the convenience of the reader, we present yet another proof of Lemma 4.2 in §5. Granted the result, we deduce from Lemma 4.1 the following.

Corollary 4.3 *If $-\Delta + V$ has discrete spectrum on $L^2(\mathbb{R}^n)$, then, for each $b \in (0, 1]$,*

$$(4.6) \quad \lambda_N^{Q_{b,\xi}}(-\Delta + V) \longrightarrow +\infty, \quad \text{as } |\xi| \rightarrow \infty.$$

Note that the left side of (4.6) is equal to

$$(4.7) \quad b^{-2} \lambda_N^Q(-\Delta + b^2 V_{b,\xi}).$$

We are now ready to prove the implication (1) \implies (3) in Theorem 1.1. Given $A \in (0, \infty)$, pick $b_0 = b_0(A)$ so small that (2.11) applies, so that

$$(4.8) \quad \Gamma(v) \leq b_0^2 A \implies \lambda_N^Q(-\Delta + v) \leq \tilde{C}_n \Gamma(v).$$

Consequently, for $b \in (0, b_0]$,

$$(4.9) \quad \Gamma(b^2 V_{b,\xi}) \leq Ab^2 \implies \lambda_N^Q(-\Delta + b^2 V_{b,\xi}) \leq \tilde{C}_n \Gamma(b^2 V_{b,\xi}) \leq \tilde{C}_n Ab^2.$$

Now, by (4.6)–(4.7), we cannot have the bound $\lambda_N^Q(-\Delta + b^2 V_{b,\xi}) \leq \tilde{C}_n Ab^2$ for large $|\xi|$, so consequently we cannot have the bound $\Gamma(b^2 V_{b,\xi}) \leq Ab^2$ for large $|\xi|$. The proof of Theorem 1.1 is complete, modulo the proof of Lemma 4.2, which will be given in the next section.

5 Proof of Lemma 4.2

Given non-negative $v \in L^1(Q)$, let us denote by L_D the operator $L = -\Delta + v$ on $L^2(Q)$ with the Dirichlet boundary condition and by L_N the operator with the Neumann boundary condition. We assume

$$(5.1) \quad \lambda = \lambda_D^Q(-\Delta + v),$$

the smallest eigenvalue of L_D , is large, and we want to estimate the smallest eigenvalue of L_N . We will estimate various “heat semigroups.” For $x, y \in Q, t > 0$, set

$$(5.2) \quad \begin{aligned} p_D(t, x, y) &= e^{-tL_D} \delta_y(x), & p_N(t, x, y) &= e^{-tL_N} \delta_y(x), \\ p_Q(t, x, y) &= e^{t\Delta_N} \delta_y(x), & p_0(t, x, y) &= (4\pi t)^{-n/2} e^{-|x-y|^2/4t}. \end{aligned}$$

Here Δ_N denotes the Laplace operator on $L^2(Q)$, with the Neumann boundary condition. It will be convenient to note the following inequalities:

$$(5.3) \quad p_D(t, x, y) \leq p_0(t, x, y), \quad p_N(t, x, y) \leq p_Q(t, x, y).$$

We want to estimate $p_N(t, x, y)$, but first we will estimate $p_D(t, x, y)$. Let us fix $a \in (0, 1)$ and set

$$(5.4) \quad \tau = \lambda^{-a}.$$

Using (5.3) we have

$$(5.5) \quad \|e^{-\tau L_D}\|_{\mathcal{L}(L^1, L^2)} \leq (4\pi\tau)^{-n/4}, \quad \|e^{-\tau L_D}\|_{\mathcal{L}(L^2, L^\infty)} \leq (4\pi\tau)^{-n/4},$$

while (5.1) gives

$$(5.6) \quad \|e^{-\tau L_D}\|_{\mathcal{L}(L^2, L^2)} \leq e^{-\tau\lambda}.$$

Hence

$$(5.7) \quad \|e^{-3\tau L_D} \delta_y\|_{L^\infty} \leq C\tau^{-n/2} e^{-\tau\lambda} = C\lambda^{an/2} e^{-\lambda^{1-a}}.$$

In other words,

$$(5.8) \quad 0 \leq p_D(3\tau, x, y) \leq C\lambda^{an/2} e^{-\lambda^{1-a}}, \quad \forall x, y \in Q.$$

Next, we estimate $V_y(t, x) = p_N(t, x, y) - p_D(t, x, y)$, for $t \in [0, 3\tau]$. We have

$$(5.9) \quad (\partial_t - L)V_y = 0 \text{ on } \mathbb{R}^+ \times Q, \quad V_y(0, x) = 0,$$

and $x \in \partial Q \implies V_y(t, x) = p_N(t, x, y)$. Hence

$$(5.10) \quad x \in \partial Q \implies 0 \leq V_y(t, x) \leq p_Q(t, x, y).$$

Let us define the set $\Omega_\tau \subset Q$ by

$$(5.11) \quad \Omega_\tau = \{y \in Q : \text{dist}(y, \partial Q) \geq \tau^{1/3}\}.$$

It is clear that, if λ is sufficiently large, so τ is sufficiently small,

$$(5.12) \quad x \in \partial Q, y \in \Omega_\tau, t \in [0, 3\tau] \Rightarrow p_Q(t, x, y) \leq Ce^{-\lambda^{a/4}},$$

so applying the maximum principle to (5.9)–(5.10) gives

$$(5.13) \quad V_y(t, x) \leq Ce^{-\lambda^{a/4}}, \quad \text{for } x \in Q, y \in \Omega_\tau, t \in [0, 3\tau],$$

and hence, by (5.8), if we take $a = 4/5$ and assume λ is sufficiently large,

$$(5.14) \quad 0 \leq p_N(3\tau, x, y) \leq C\lambda^{2n/5}e^{-\lambda^{1/5}}, \quad \forall x \in Q, y \in \Omega_\tau.$$

Now, using the semigroup property of e^{tL_N} and the fact that $\|e^{-tL_N}\|_{\mathcal{L}(L^\infty, L^\infty)} \leq 1$, we deduce that

$$(5.15) \quad 0 \leq p_N(t, x, y) \leq C\lambda^{2n/5}e^{-\lambda^{1/5}}, \quad \forall x \in Q, y \in \Omega_\tau, t \geq 3\tau.$$

In particular, if λ is large enough that $3\tau = 3\lambda^{-4/5} < 1$, the estimate (5.15) applies with $t = 1$. On the other hand, we can use (5.3) to obtain

$$(5.16) \quad p_N(1, x, y) \leq p_Q(1, x, y) \leq C, \quad \forall x \in Q, y \in Q \setminus \Omega_\tau.$$

It follows that

$$(5.17) \quad \begin{aligned} \int_Q p_N(1, x, y) dy &\leq C\lambda^{2n/5}e^{-\lambda^{1/5}} + C \text{Vol}(Q \setminus \Omega_\tau) \\ &\leq C\lambda^{2n/5}e^{-\lambda^{1/5}} + C\lambda^{-4/15}. \end{aligned}$$

Of course $p_N(1, x, y) = p_N(1, y, x)$, so there is a similar bound on $\int_Q p_N(1, x, y) dx$. Hence we deduce that

$$(5.18) \quad \|e^{-L_N}\|_{\mathcal{L}(L^2, L^2)} \leq C\lambda^{2n/5}e^{-\lambda^{1/5}} + C\lambda^{-4/15} = \Phi(\lambda).$$

It follows that

$$(5.19) \quad \lambda_N^Q(-\Delta + \nu) \geq \log \frac{1}{\Phi(\lambda)},$$

and Lemma 4.2 is proven.

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