

**On an Asymptotic Expansion of the Hypergeometric Function.**

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In a previous paper\* the author has employed the expansion

$$F(\alpha, \beta, \gamma, z) = \sum_{r=1}^s T_r + P_s, \dots\dots\dots(1)$$

where

$$T_r = \frac{\alpha(\alpha+1)\dots(\alpha+r-2)\beta(\beta+1)\dots(\beta+r-2)}{(r-1)! \gamma(\gamma+1)\dots(\gamma+r-2)} z^{r-1},$$

$$P_s = \frac{1}{B(\beta+s, \gamma-\beta)} T_{s+1} \int_0^1 s(1-t)^{s-1} I dt,$$

and † 
$$I = \int_0^1 \xi^{\beta+t-1} (1-\xi)^{\gamma-\beta-1} (1-\xi tz)^{-\alpha-s} d\xi,$$

to establish the theorem that, if  $-\pi/2 < \text{amp } \gamma < \pi/2$ , the function  $P_s/T_{s+1}$  remains finite as  $\gamma \rightarrow \infty$ . This theorem is valid provided that  $z$  is not real and  $\geq 1$ .

It will now be shown that the theorem is true for a more extended range of values of  $\text{amp } \gamma$ .

In the integral  $I$  put  $\zeta = 1 - e^{-\lambda} \ddagger$ ; then

$$I = \int_0^\infty (1 - e^{-\lambda})^{\beta+s-1} e^{-\lambda(\gamma-\beta)} \{1 - tz(1 - e^{-\lambda})\}^{-\alpha-s} d\lambda, \dots\dots\dots(2)$$

the path of integration being the real axis from 0 to  $\infty$ . This path may now be deformed into that straight line from the origin to infinity which makes an acute angle  $-\phi$  with the positive real axis, and the integral is still convergent provided that, on the path of integration,  $1 - tz(1 - e^{-\lambda}) \neq 0$ , and provided that

$$-\pi/2 < \text{amp } (\lambda\gamma) < \pi/2.$$

\* *Proc. Edin. Math. Soc.*, Vol. XLI., pp. 82-92.

†  $R(\gamma)$  and  $s$  are taken so large that  $R(\gamma - \beta) > 0$  and  $R(\beta + s) > 0$ .

‡ Cf. G. N. Watson, *Trans. Camb. Phil. Soc.*, Vol. 22, 1918, p. 299.

Since  $\text{amp } \lambda = -\phi$  the latter inequality may be written

$$-\pi/2 + \phi < \text{amp } \gamma < \pi/2 + \phi,$$

and, by reversing the transformation  $\zeta = 1 - e^{-\lambda}$ , it can be made clear that the first condition may be replaced by the proviso that  $1 - tz\zeta$  must not vanish at any point on the contour in the  $\zeta$ -plane which corresponds to the path of integration in the  $\lambda$ -plane.

To determine this contour put  $\lambda = \mu(1 - it)$ , where  $\mu$  is real and  $t = \tan\phi$ ; then

$$\zeta = \xi + i\eta = 1 - e^{-\lambda} = 1 - e^{-\mu + i\mu t},$$

so that

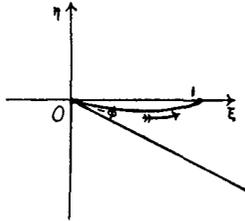
$$\xi = 1 - e^{-\mu} \cos \mu t, \quad \eta = -e^{-\mu} \sin \mu t.$$

Hence, if  $(r, \theta)$  are the polar coordinates of the point  $(\xi, \eta)$  referred to axes parallel to the  $\xi$  and  $\eta$  axes and passing through the point  $(1, 0)$  in the  $\zeta$ -plane

$$r = e^{-\mu}, \quad \tan \theta = \tan \mu t,$$

and the contour is an equiangular spiral (see Fig.) whose equation may be written

$$r = e^{-(\theta + \pi) \cot \phi},$$



where  $\theta = -\pi$  and  $r = 1$  when  $\mu = 0$ ,  $\lambda = 0$ ,  $\zeta = 0$ ; at this point  $\frac{d\zeta}{d\lambda} = 1$ , so that the contour makes an angle  $-\phi$  with the  $\xi$ -axis, and the path is described from  $\zeta = 0$  in the direction indicated by the arrow.

It can easily be shown that the entire contour of integration lies between the lines  $\text{amp } \zeta = \pm \phi$ . Now  $1 - tz\zeta$  must not vanish for any point  $\zeta$  on this contour. But, since  $0 \leq t \leq 1$ , the values of  $z$  which satisfy  $z = 1/(t\zeta)$  will lie in the region to the right of the  $\eta$ -axis which is bounded by the lines  $\text{amp } \zeta = \pm \phi$ . Also the

hypergeometric expansion is valid within the unit circle; hence the expansion (1) is valid for

$$-\pi/2 + \phi < \text{amp } \gamma < \pi/2 + \phi$$

at all points external to a region  $B$  which is bounded by the lines  $\text{amp } \zeta = \pm \phi$  and the circle  $|\zeta| = 1$ .

But if  $z$  is any interior point of the region  $A$  consisting of the entire complex plane bounded by a cross-cut along the positive real axis from  $+1$  to  $+\infty$ ,  $\phi$  can be chosen so small ( $< |\text{amp } z|$ ) that  $z$  does not lie in  $B$ . Hence for any interior point of  $A$  a  $\phi$  can be found such that the expansion defined by (1) and (2) is valid for  $-\pi/2 + \phi < \text{amp } \gamma < \pi/2 + \phi$ .

It remains to prove that the theorem is true under these conditions. Now, for points on the path of integration in the  $\lambda$ -plane

$$\begin{aligned} |1 - e^{-\lambda}| &= \sqrt{(1 - 2e^{-\mu} \cos \mu t + e^{-2\mu})} \\ &= \sqrt{\{(1 - e^{-\mu})^2 + 2e^{-\mu}(1 - \cos \mu t)\}} \\ &= (1 - e^{-\mu}) \sqrt{\left\{1 + \left(\frac{\sin \frac{1}{2} \mu t}{\sinh \frac{1}{2} \mu}\right)^2\right\}} \\ &< C(1 - e^{-\mu}), \end{aligned}$$

where  $C$  is a definite positive constant ( $|\phi| < \pi/2$ ). Also, let  $\gamma = ge^{i\phi}$ , so that  $-\pi/2 < \text{amp } g < \pi/2$ , and note that  $\lambda = \mu \sec \phi e^{-i\phi}$ . Then, if  $\beta = \sigma + i\tau$  and  $g = l + im$ ,

$$R\{\lambda(\gamma - \beta)\} = \mu(l \sec \phi - \sigma - \tau \tan \phi).$$

Accordingly, for any point  $z$  within the region  $A$ ,

$$|I| < K \int_0^\infty (1 - e^{-\mu})^{\sigma+s-1} e^{-\mu(l \sec \phi - \sigma - \tau \tan \phi)} d\mu,$$

where  $K$  is a definite constant. Here put  $\xi = 1 - e^{-\mu}$  and get

$$\begin{aligned} |I| &< K \int_0^1 \xi^{\sigma+s-1} (1 - \xi)^{l \sec \phi - \sigma - \tau \tan \phi - 1} d\xi \\ &= KB(\sigma + s, l \sec \phi - \sigma - \tau \tan \phi); \end{aligned}$$

thus

$$\begin{aligned} \left| \frac{P_s}{T_{s+1}} \right| &< K \frac{B(\sigma + s, l \sec \phi - \sigma - \tau \tan \phi)}{|B(\beta + s, \gamma - \beta)|} \\ &= K \frac{\Gamma(\sigma + s)}{|\Gamma(\beta + s)|} \frac{\Gamma(l \sec \phi - \sigma - \tau \tan \phi)}{|\Gamma(l \sec \phi - \tau \tan \phi + s)|} \left| \frac{\Gamma(\gamma + s)}{\Gamma(\gamma - \beta)} \right|, \end{aligned}$$

and when  $\gamma \rightarrow \infty$ ,  $l \rightarrow \infty$ , and

$$\left| \frac{P_s}{T_{s+1}} \right| < K \frac{\Gamma(\sigma + s)}{|\Gamma(\beta + s)|} \frac{|\gamma^{\beta+s}|}{(l \sec \phi)^{\sigma+s}} = K \frac{\Gamma(\sigma + s)}{|\Gamma(\beta + s)|} \left( \frac{|\gamma|}{l \sec \phi} \right)^{\sigma+s} e^{-\chi\tau},$$

where  $\text{amp } \gamma = \chi$ . But  $|\gamma|/l$  is finite; hence the theorem holds.

Similarly the theorem can be shown to hold for the region  $-\pi/2 - \phi < \text{amp } \gamma < \pi/2 - \phi$ ; thus it holds for the entire region

$$-\pi/2 - \phi < \text{amp } \gamma < \pi/2 + \phi.$$

*The Asymptotic Expansion of  $P_n^m(z)$ .* It follows that, for any interior point  $z$  of the region in which the asymptotic expansion of  $P_n^m(z)$  for  $n$  large is valid, a  $\phi$  can be found such that the asymptotic expansion holds for  $-\pi/2 - \phi < \text{amp } n < \pi/2 + \phi$ . Hence, as  $P_{-n-1}^m(z) = P_n^m(z)$ , the function possesses an asymptotic expansion for any value of  $\text{amp } n$ .

By means of the formula

$$Q_n^m(z) = Q_{-n-1}^m(z) + \frac{\pi e^{m\pi i} \cos n\pi}{\sin(n-m)\pi} \frac{\Pi(n+m)}{\Pi(n-m)} P_{-n-1}^{-m}(z)$$

a similar result can be obtained for the function  $Q_n^m(z)$ .

