

# Ultrashort pulse capability at the L2I high intensity laser facility

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## Abstract

The Laboratory for Intense Lasers (L2I) is a research centre in optics and lasers dedicated to experimental research in high intensity laser science and technology and laser plasma interaction. Currently the laboratory is undergoing an upgrade with the goal of increasing the versatility of the laser systems available to the users, as well as increasing the pulse repetition rate. In this paper we review the current status of the laser research and development programme of this facility, namely the upgraded capability and the recent progress towards the installation of an ultrashort, diode-pumped OPCPA laser system.

**Keywords:** diode-pumped solid-state laser and applications; femtosecond laser and applications; frequency conversion; laser amplifiers; nonlinear optics

## 1. Introduction

The Laboratory for Intense Lasers (L2I) is a high intensity laser facility hosted by Instituto Superior Técnico at the University of Lisbon. It is dedicated to experimental research in high intensity laser science and technology, ultrashort diagnostics and laser plasma interaction, with emphasis in plasma particle accelerators, high harmonic generation and advanced radiation sources. L2I has been recently selected to a top group of research infrastructures of national strategic relevance, in the scope of which it is undergoing a renovation programme with the goal of enabling external users to benefit from the available and future laser systems.

The main laser of L2I is a 15 TW, hybrid Ti:sapphire–Nd:glass system, based on the chirped pulse amplification technique and operating at 1053 nm. This laser has been serving internal users for several years, e.g., for the development and testing of plasma diagnostics<sup>[1]</sup>, guiding schemes for plasma accelerators<sup>[2–4]</sup> and high harmonic generation<sup>[5]</sup>.

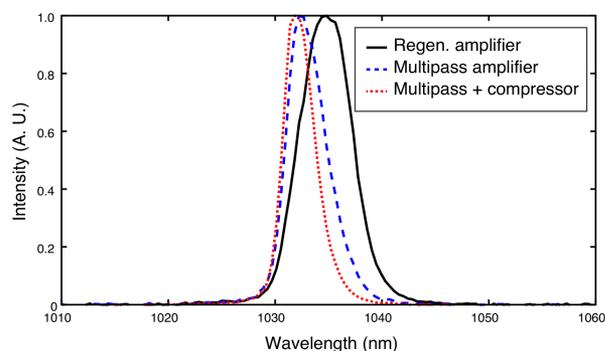
Over the past years we have been developing a diode-pumped laser programme, with the goal of improving the

experimental capability of the laboratory, namely by improving the repetition rate from the mHz to the Hz rate. We have successfully concluded a dual-stage, diode-pumped amplifier based on Yb:CaF<sub>2</sub> and Yb:YAG operating at 1030 nm and delivering 100 mJ, sub-ps pulses at 1 Hz, and made this laser available for high harmonic generation experiments. Simultaneously this amplifier has been designed to pump an ultrabroadband optical parametric chirped pulse amplifier (OPCPA) based on the nonlinear crystal yttrium calcium oxyborate (YCOB), with the goal of installing a homemade 20 mJ, sub-20 fs laser source. This crystal has been shown to exhibit the required qualities for J-level, few-fs amplification based in OPCPA<sup>[6–9]</sup> and has the potential to become an efficient alternative to current options such as BBO, LBO or (D)KDP for high peak and high average power amplification.

In this paper we describe the current status of the implementation of high-power, ultrashort pulse capability at L2I, in particular evaluating the potential of the ultrabroadband OPCPA system. We start by describing the design and performance of the 100-mJ-level diode-pumped amplifier and the creation of a broadband, low energy signal by white-light continuum generation in bulk media. Next we describe the concept for ultrabroadband noncollinear OPCPA in YCOB and evaluate recent results and perspectives.

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**Figure 2.** Output spectrum from regenerative amplifier (black), multipass amplifier (blue) and grating compressor (red).

into the multipass amplifier. However, in the OPCPA configuration a fraction ( $\sim 1.2$  mJ) of this output is compressed, either using the same CVBG or the grating compressor, resulting in a compressed pulse duration of 400 fs in the first case and 230 fs in the second. This compressed mJ-level pulse is then focused on a 10 mm thick sapphire block, and through white-light continuum (WLC) generation an ultrabroadband (650–1000 nm) signal for seeding the OPCPA is produced. Results for WLC generated with a similar setup can be found in Refs. [11, 12].

The Yb:YAG multipass amplifier is currently capable of delivering a maximum output energy of 106 mJ at 1 Hz, limited by optical damage to the final mirror surface, for a pump power of 1.4 kW. The output spectrum peak is shifted to 1032 nm, with an FWHM spectral bandwidth of  $\sim 4$  nm. When used for experiments, the amplified pulses are sent to the grating pair compressor mentioned earlier, resulting in a compressed pulse duration of 730 fs, with a compressor efficiency of 54%. Pulse spectra were acquired by using a high resolution, broadband spectrometer (Ocean Optics HR2000+) coupled to an integration sphere (Ocean Optics model FOIS-1). Figure 2 shows the evolution of the spectrum transmitted by the two amplifiers and by the compressor, evidencing the mentioned spectral shift between the amplifiers. This shift can be explained by the different spectral properties of Yb:CaF<sub>2</sub> and Yb:YAG<sup>[13]</sup>. In the first case, the balance between a smoothly decreasing gain cross section and an also decreasing absorption between 1030 and 1040 nm results in an effective amplification peak at  $\sim 1034.5$  nm. When injected with the broadband signal from the oscillator, this central wavelength is preserved (black curve in Figure 2). In the second case, Yb:YAG exhibits a very sharp absorption peak at  $\sim 1032$  nm, coincident with a local, broad absorption region, resulting in a strong, narrowband gain around this wavelength. This leads to the pronounced gain-induced shift and slight gain-induced narrowing (5.7–4.0 nm) that is observed.

Figure 3 shows an image of the beam profile (acquired using an intensity and wavefront meter, Phasics model SID4)

evidencing the excellent spatial quality of the amplifier and compressor. In the OPCPA pumping configuration, the uncompressed output is directly frequency doubled in a 15 mm long DKDP crystal, resulting in 40 mJ pulses centred at 516 nm. These are then directed to the double-stage YCOB amplifier, as described next.

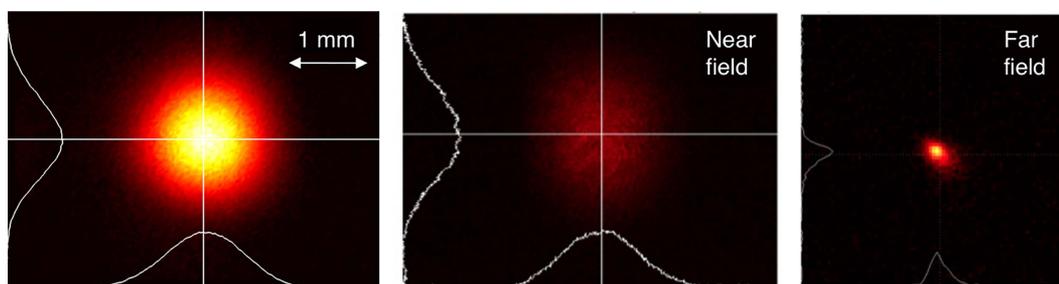
### 3. Noncollinear, ultrabroadband OPCPA based on YCOB

In general, a noncollinear configuration for optical parametric amplification allows the broadest phase matching bandwidth<sup>[14]</sup>. For these conditions, YCOB has been shown to offer a very large gain bandwidth, in excess of 130 nm<sup>[9]</sup>, which is remarkable even when compared to BBO or LBO. It is also known that for YCOB, being a biaxial crystal, the highest value of the nonlinear coefficient is obtained using a combination of phase matching angles outside of the principal planes<sup>[6, 15]</sup>. Other advantages include its large transparency range and high damage threshold. In addition, recent progress in the growth of YCOB has made available 100 mm aperture boules, showing the great potential of this nonlinear medium for high energy, ultrahigh intensity optical parametric amplification.

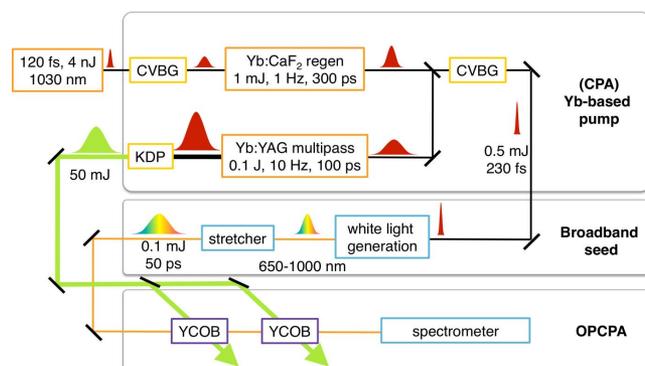
We have developed a double-stage noncollinear OPCPA based in YCOB crystals cut with their angles optimized for maximum nonlinear coefficient. Figure 4 shows the schematic of the overall setup. Pulses from the modelocked oscillator are first amplified in the two-stage diode-pumped amplifier described earlier. The output of the regenerative amplifier is split into two arms by using an 80/20 beam-splitter. As mentioned above, the major fraction of the energy is compressed in a CVBG, and focused in a sapphire crystal, leading to white-light continuum generation. A suitable bandpass filter limits the continuum to the range 650–1000 nm, which is used as the seed for the parametric amplifier. The remaining fraction of the regenerative amplifier is sent to the multipass amplifier, amplified to the 100 mJ level and frequency doubled in a KDP crystal. The resulting 50 mJ of green light is then split and used to pump both YCOB crystals.

We started by developing a 1D numerical model for OPCPA in a general biaxial crystal, and used the parameters for YCOB in order to evaluate the gain and phase matching bandwidth for an amplifier pumped at 515 nm<sup>[6]</sup>. This has resulted in an optimum noncollinear angle of  $3.5^\circ$  and optimum phase matching angles  $\theta = 115^\circ$  and  $\phi = -43.5^\circ$  for a 15 mm crystal. For a signal centred at 800 nm, this yields a (continuous FWHM) maximum gain bandwidth of 110 nm.

As mentioned earlier, the seed pulses for the OPCPA are generated by taking the mJ-level output of the regenerative amplifier and generating a WLC in a sapphire block. The



**Figure 3.** Near-field beam profile at the output of the multipass amplifier (left); near (centre) and far (right) beam profile at the grating compressor output.

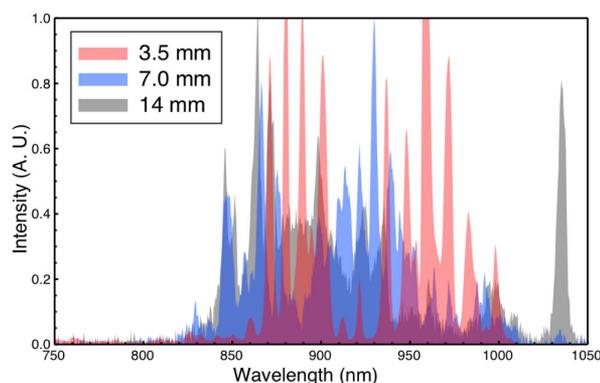


**Figure 4.** Schematic of diode-pumped OPCPA system.

great advantage of this approach is that it ensures optical synchronization between pump and signal, and they remain separated only for a distance corresponding to the optical path of the multipass amplifier ( $\sim 9$  m).

The continuum was characterized by a homemade cross correlation frequency resolved optical gating (XFROG) device in order to evaluate the distribution of frequencies over time and the operating zone for amplification and compression. To overcome phase matching limitations, crystal dithering must be used<sup>[16]</sup>. The full spectrum was then retrieved using commercial software, resulting in a usable region extending from 650 to 1000 nm<sup>[11]</sup>. One advantage of using a pump at 1030 nm is that the original pulse can be removed by using a suitable high pass filter. One limitation of this method is that the onset of multiple filamentation due to crystal damage may sometimes prevent obtaining a smooth and regular spectrum. Typically the generated white light contains a central white portion where the supercontinuum is stable, and a blue-green outer ring that is removed by an iris.

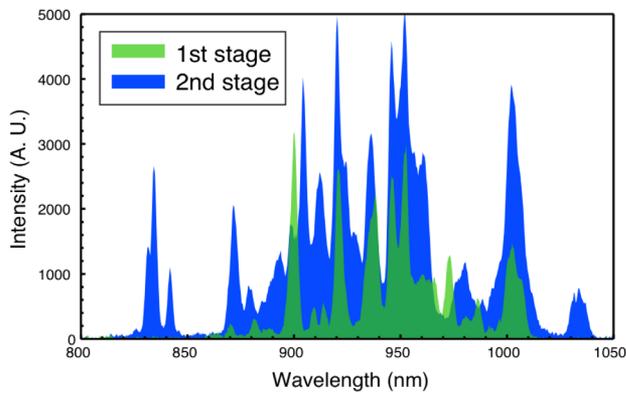
The WLC is stretched to  $\sim 50$  ps and chirped in a 125 mm long block of SF11 glass. The pulses are then time-delayed in order to account for the optical path of the multipass amplifier and the frequency doubling stage of the pump. We started by comparing the amplification performance for three YCOB crystals (3.5, 7 and 14 mm long) at a  $\sim 3.5^\circ$  noncollinear angle, at the first amplifier stage. This first



**Figure 5.** Maximum optimum amplification bandwidth versus YCOB crystal length: 3.5 mm (pink), 7 mm (blue) and 14 mm (grey).

amplifier is run at very high gain, with no considerable saturation or reconversion, due to the very low energy of the signal pulses ( $< 5 \mu\text{J}$ ). Since pump depletion was negligible even for the longest crystal, we decided to use the full pump power on the first stage and to recycle it for the second crystal. The maximum output spectral bandwidth obtained for each of the crystals is shown in Figure 5. For each case, a small adjustment of the angles can be made in order to either obtain a narrower continuous bandwidth, or an equally long but interrupted amplification bandwidth. As expected, the longer crystals enable a larger, less modulated bandwidth. Both the 7 and 14 mm long crystals could provide up to 150 nm full width at  $1/e^2$ . For the longer crystal it was also possible to amplify a continuous bandwidth encompassing up to 1030 nm by using a slightly larger noncollinear angle. The 3.5 mm crystal could exhibit gain over 100 nm, although this occurred mostly at the peak of the supercontinuum leading to a highly modulated spectrum.

We performed a preliminary evaluation of two-stage amplification, but due to space limitations it was not possible to increase the pump diameter sufficiently in order to avoid optical damage after the first amplification stage. By carefully tuning the phase matching angle of the second OPCPA stage we were able to increase the output bandwidth by slightly detuning the gain away from the pulse central wavelength towards  $\sim 920$  nm (Figure 6). The resulting amplified



**Figure 6.** Gain bandwidth for the first (green) and second (blue) OPCPA stages.

bandwidth spanned over 225 nm (of which 150 nm were continuous and above the  $1/e^2$  level). Such bandwidth is enough to support pulse compression to the few-fs level, well within the goals of generating sub-20 fs pulses. Currently we are working on the implementation of the compression stage, based on transmission gratings, prisms and chirped mirrors.

#### 4. Conclusions

We have described the implementation of high energy, ultrashort pulse capability at L2I, following the progress of our diode-pumped programme. We have successfully developed and commissioned a 100 mJ, 1 Hz diode-pumped amplifier based on Yb:CaF<sub>2</sub> and Yb:YAG operating at 1030 nm. This amplifier can work as an independent high-power pulse source, by having the output pulses compressed to the sub-ps level, resulting in  $\sim 0.1$  GW pulses at 1 Hz. Alternatively the output of the amplifier can be frequency doubled and used to pump an optically synchronized, two-stage OPCPA based on the nonlinear crystal YCOB. The tests performed so far have evidenced the ultrabroadband nature of this material, resulting in amplified gain bandwidths in excess of 200 nm, corresponding to few-cycle Fourier transform limited pulse durations. Future work will include the installation and characterization of an ultrabroadband compressor based on a combination of transmission gratings, prisms and specially designed chirped mirrors. Once completed, the new ultrashort laser system will be made available to users.

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