## ON A THEOREM IN THE GENERALISED FOURIER TRANSFORM

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1. Introduction . The function  $\omega_{\mu, v}$ (x) was defined by G.N. Watson, [9, (i)] in 1931 by the integral relation \*

$$\widetilde{\omega}_{\mu, v}(x) = x^{\frac{1}{2}} \int_{0}^{\infty} J_{\mu}(t) J_{v}(x/t) t^{-1} dt$$

$$= \frac{x^{v+\frac{1}{2}} 2^{-2v-1} \Gamma(\frac{1}{2}\mu - \frac{1}{2}v)}{\Gamma(v+1)\Gamma(1 + \frac{1}{2}\mu + \frac{1}{2}v)} {}_{0}F_{3}(v+1, 1 - \frac{1}{2}\mu + \frac{1}{2}v, 1 + \frac{1}{2}\mu + \frac{1}{2}v; x^{2}/16)$$

+ another term with  $\mu$  and v interchanged;

$$-R(\mu + \frac{3}{2}) < 0 < R(v + \frac{3}{2}).$$

He showed (without proof) that it is a symmetric Fourier kernel. Later K.P. Bhatnagar, [1, (i), (ii)] in 1953 and 1954 investigated in some details the properties of this kernel and extended it to n parameters and defined

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<sup>\*</sup> The integral  $\int_{0}^{\infty} J_{\mu}(t)J_{\nu}(\frac{x}{t}) t^{\rho} dt$  was originally evaluated by C.V.H. Rao. See Messenger of Maths, 47, (1918), 134-7. Also see Bessel Functions by Watson. (1922) 437.

$$\widetilde{\omega}_{\mu_{1}}, \dots, \mu_{n}(x) = x^{\frac{1}{2}} \int_{0}^{\infty} \dots \int_{0}^{\infty} J_{\mu_{1}}(t_{1}) \dots J_{\mu_{n-1}}(t_{n-1}) J_{\mu_{n}}(\frac{x}{t_{1} \dots t_{n-1}}) \times (t_{1} \dots t_{n-1})^{-1} dt_{1} \dots dt_{n-1}$$

$$= \int_{0}^{\infty} \widetilde{\omega}_{\mu_{1}}, \dots \mu_{n-1}(x/t) J_{\mu_{n}}(t)^{-\frac{1}{2}} dt,$$

where  $R(\mu_k + \frac{1}{2}) \ge 0$ , k = 1, 2, ..., n and the  $\mu$ 's may be permuted among themselves.

V.P. Mainra, [3, (ii)] in 1958 defined the function

$$\widetilde{\omega}_{\mu, v}^{\lambda}(x) = \int_{0}^{\infty} \widetilde{\omega}_{\mu, v}(xt) J_{\lambda}(t) \sqrt{t} dt,$$

and further generalised this kernel as

$$\omega_{\mu_{1},\ldots\mu_{n}}^{v_{1},\ldots v_{m}}(\mathbf{x}) = \int_{0}^{\infty} \ldots \int_{0}^{\infty} J_{v_{1}}(t_{1}) \ldots J_{v_{m}}(t_{m}) J_{\mu_{1}}(T_{1}) \ldots J_{\mu_{n-2}}(T_{n-2}) \times U_{n}^{v_{1},\ldots v_{m}}(\mathbf{x}) = \int_{0}^{\infty} \ldots \int_{0}^{\infty} J_{v_{1}}(t_{1}) \ldots J_{v_{m}}(t_{m}) J_{\mu_{1}}(T_{1}) \ldots J_{\mu_{n-2}}(T_{n-2}) \times U_{n}^{v_{1},\ldots v_{m}}(\mathbf{x}) = \int_{0}^{\infty} \ldots \int_{0}^{\infty} J_{v_{1}}(t_{1}) \ldots J_{v_{m}}(t_{m}) J_{\mu_{1}}(T_{1}) \ldots J_{\mu_{n-2}}(T_{n-2}) \times U_{n}^{v_{1},\ldots v_{m}}(t_{n}) + U_{n}^{v_{1},\ldots v_{m}}(t_{n}) J_{\mu_{1}}(T_{1}) \ldots J_{\mu_{n-2}}(T_{n-2}) \times U_{n}^{v_{1},\ldots v_{m}}(t_{n}) + U_{n}^{v_{1},\ldots v_{m}}($$

$$\widetilde{\omega}_{\mu_{n-1},\mu_n}(\underbrace{\frac{xt_1\dots t_m}{T_1\dots T_{n-2}}}_{1}) \cdot \underbrace{\frac{(t_1\dots t_m)^{\frac{1}{2}}}{T_1\dots T_{n-2}}}_{1} \, \mathrm{d}T_1\dots \mathrm{d}T_{n-2} \, \mathrm{d}t_1\dots \mathrm{d}t_m \;,\; m < n.$$

$$R(\mu_r, v_s) > -1,$$
  $r = 1, 2, ...n,$   $s = 1, ...m.$ 

In 1961, Charles Fox [2] showed that the G-Function as defined by C.S. Meijer, [4] is a symmetric Fourier kernel. For certain values of the parameters the kernel degenerates into the kernels of Bhatnagar and Mainra, but he did not investigate the properties of this kernel. B. Singh [5, (i)] in 1964, Y.P. Singh [7, (i)] and P. Singh [6] in 1966 showed that many of the results given in Watson's Bessel Functions follow as particular cases of the integrals involving the generalised kernels mentioned above.

In the present paper, the author has proved certain interesting results involving the functions f(t), g(t) and their transforms F(x) and G(x) in these generalised transforms.

The following results are either known or can be proved easily.

(1) 
$$\omega_{\mu}(x) = \sqrt{x} J_{\mu}(x)$$
,  $\omega_{\mu, \mu+1}(x) = J_{2\mu+1}(2\sqrt{x})$ ,  $R(\mu) > -1$ .

(2) 
$$\tilde{\omega}_{\mu_{1}, \dots \mu_{n}}(x) = \int_{0}^{\infty} \tilde{\omega}_{\mu_{1}, \dots \mu_{n-1}}(xt) J_{\mu_{n}}(t^{-1})t^{-\frac{3}{2}} dt$$

$$= \int_{0}^{\infty} \tilde{\omega}_{\mu_{1}, \dots \mu_{r}}(xt) \tilde{\omega}_{\mu_{n+1}, \dots \mu_{n}}(t^{-1}) \frac{dt}{t}, R(\mu_{r}) > -1.$$

$$r = 1, 2, \dots n.$$

(4) 
$$\omega_{\mu_{1}, \dots \mu_{n}}^{(x)} = 0 \left(x^{\frac{\mu_{r} + \frac{1}{2}}{2}}\right), \quad r = 1, 2, \dots n \text{ for small } x$$

$$= \frac{\frac{1-n}{2n}}{x^{2n}} \left[\cos \left(2nx^{\frac{n}{n}} + \alpha\right)(A + 0(x^{\frac{n}{n}}))\right]$$

$$+ \sin(2nx^{\frac{1}{n}} + \alpha) 0(x^{\frac{1}{n}})$$

$$= \frac{\frac{1-n+m}{2(n-m)}}{x^{2}(n-m)} \left[\cos((2n-m)x^{\frac{1}{n-m}} + \alpha)(A_{1} + 0(x^{\frac{n-m}{n-m}})\right]$$

+ 
$$\sin(2(n-m) \times \frac{1}{n-m} + \alpha_1) 0(x^{n-m}) + \sum_{n=1}^{m} x^{-2(\frac{3}{4} + \frac{a_n}{2})} \times$$

$$\{P_n + 0(x^{-2})\}$$
,

for large x and  $\alpha$ ,  $\alpha_1$ , A, A, P are constants, =  $0(x^{\mu/2} + \frac{1}{2})$  for small x, r = 1, 2, ...n.

(6) The Mellin transforms of  $\tilde{\omega}_{\mu_1}, \dots, \mu_n$  (x) and

and

$$2^{(n-m)(s-\frac{1}{2})} \frac{\Gamma(\frac{\mu_{1}}{2}+\frac{s}{2}+\frac{1}{4})\dots\Gamma(\frac{\mu_{n}}{2}+\frac{s}{2}+\frac{1}{4})\Gamma(\frac{v_{1}}{2}-\frac{s}{2}+\frac{3}{4})\dots\Gamma(\frac{v_{m}}{2}-\frac{s}{2}+\frac{3}{4})}{\Gamma(\frac{\mu_{1}}{2}-\frac{s}{2}+\frac{3}{4})\dots\Gamma(\frac{\mu_{n}}{2}-\frac{s}{2}+\frac{3}{4})\Gamma(\frac{v_{1}}{2}+\frac{s}{2}+\frac{1}{4})\dots\Gamma(\frac{v_{m}}{2}+\frac{s}{2}+\frac{1}{4})}$$

and belong to L(- $\infty$ ,  $\infty$ ) if  $\frac{1}{2} - \frac{1}{n} > \sigma$ ,  $s = \sigma + it$ .

Results (2) and (3) can be proved by an application of Parseval's theorem [8, p. 54].

## Notations employed

$$f_{\mu_1,\ldots,\mu_n}(x) = \int_0^\infty f(t) \tilde{\omega}_{\mu_1,\ldots,\mu_n}(xt) dt$$

$$f_{\mu_4,\ldots,\mu_n}^{v_1,\ldots v_m}(x) = \int_0^\infty f(t) \omega_{\mu_4,\ldots,\mu_n}^{v_1,\ldots v_m}(xt) dt$$

If  $g(x) = \int_{0}^{\infty} f(t) \tilde{\omega}_{\mu_{1}, \dots, \mu_{n}}(xt) dt$ , g(x) is called the  $\tilde{\omega}_{\mu_{1}, \dots, \mu_{n}}(x)$  transform of f(x).

If 
$$g(x) = f(x)$$
,  $f(x)$  is said to be  $R_{\mu_1}, \dots \mu_n$ .

THEOREM 1. Let f(t) and G(t) be continuous and belong to  $L(0,\infty)$  and F(x), G(x) be the  $\omega$   $\mu_1,\ldots\mu_n$ (x) transforms of f(t) and g(t) respectively. Then

$$\int_{0}^{\infty} F(x) G(x) dx = \int_{0}^{\infty} f(t)g(t) dt.$$

Proof.

$$\int_{0}^{\infty} F(x)G(x)dx = \int_{0}^{\infty} G(x)dx \int_{0}^{\infty} f(t) \tilde{\omega}_{\mu_{1}}, \dots \mu_{n}^{(xt)dt}$$

$$= \int_{0}^{\infty} f(t)dt \int_{0}^{\infty} G(x)\tilde{\omega}_{\mu_{1}}, \dots \mu_{n}^{(xt)dt}$$

$$= \int_{0}^{\infty} f(t)g(t)dt.$$

This is the Parseval theorem for the transform  $\overset{\sim}{\omega}_{\mu_1}, \ldots \overset{\sim}{\mu}_n$  introduced by Bhatnagar [1, (i)].

THEOREM 2. Let f(t), G(t) and g(t) satisfy the conditions of Theorem 1. Then

$$\int_{0}^{\infty} \mathbf{F}(\mathbf{x}) \mathbf{G}(\mathbf{x}) d\mathbf{x} = \int_{0}^{\infty} f_{\mu_{1}}(t) \mathbf{g}_{\mu_{1}}(t) dt$$

$$= \dots = \int_{0}^{\infty} f_{\mu_{n}}(t) \mathbf{g}_{\mu_{n}}(t) dt .$$

Proof.

$$\int_{0}^{\infty} f(t) \widetilde{\omega}_{\mu_{1}}, \dots \mu_{n}$$
 (xt) dt
$$= \int_{0}^{\infty} \sqrt{t} f(t) dt \int_{0}^{\infty} \widetilde{\omega}_{\mu_{1}}, \dots \mu_{r-1}, \mu_{r+1} \dots \mu_{n}^{(xy)} J_{\mu_{r}} (t/y) y^{-\frac{3}{2}} dy,$$

$$R(\mu_{s}) > -1, \quad s = 1, 2, \dots n,$$

$$= \int_{0}^{\infty} f_{\mu_{r}} (y^{-1}) \widetilde{\omega}_{\mu_{1}}, \dots \mu_{r-1}, \dots \mu_{n}^{(xy)} y^{-1} dy .$$

Therefore

$$\int_{0}^{\infty} G(x) dx \int_{0}^{\infty} f(t) \widetilde{\omega}_{\mu_{1}}^{\alpha}, \dots \mu_{n}^{\alpha} (xt) dt$$

$$= \int_{0}^{\infty} G(x) dx \int_{0}^{\infty} f_{\mu_{r}}^{\alpha} (y^{-1}) \widetilde{\omega}_{\mu_{1}}^{\alpha}, \dots \mu_{r-1}^{\alpha}, \mu_{r+1}^{\alpha}, \dots \mu_{n}^{\alpha} (xy) dy/y$$

$$= \int_{0}^{\infty} f_{\mu_{r}}^{\alpha} (y^{-1}) y^{-1} dy \int_{0}^{\infty} G(x) \widetilde{\omega}_{\mu_{1}}^{\alpha}, \dots \mu_{r-1}^{\alpha}, \mu_{r+1}^{\alpha}, \dots \mu_{n}^{\alpha} (xy) dx.$$

Now

$$\int_{0}^{\infty} G(x) \widetilde{\omega}_{\mu_{1}, \dots \mu_{r-1}, \mu_{r+1}, \dots \mu_{n}}^{(x) dx}$$

$$= \int_{0}^{\infty} \widetilde{\omega}_{\mu_{1}, \dots \mu_{r-1}, \mu_{r+1}, \dots \mu_{n}}^{(xy) dx} \int_{0}^{\infty} g(t) \widetilde{\omega}_{\mu_{1}, \dots \mu_{n}}^{(xt) dt}$$

$$= \int_{0}^{\infty} g(t) dt \int_{0}^{\infty} \widetilde{\omega}_{\mu_{1}, \dots \mu_{r-1}, \mu_{r+1}, \dots \mu_{n}}^{(x) \omega} (x^{x}) \widetilde{\omega}_{\mu_{1}, \dots \mu_{n}}^{(xt) dt}$$

$$= y^{-1} \int_{0}^{\infty} g(t) \omega_{\mu_{1}, \dots \mu_{r}}^{\mu_{1}, \dots \mu_{r-1}, \mu_{r+1}, \dots \mu_{n}} (t/y) dt$$

$$= y^{-1} \int_{0}^{\infty} g(t) \omega_{\mu_{r}}^{(t/y)} (t/y) dt = y^{-1} \int_{0}^{\infty} g(t) \sqrt{\frac{t}{y}} J_{\mu_{r}}^{(t/y)} dt$$

$$= y^{-1} g_{\mu} (y^{-1}) dy.$$

Therefore

$$\int_{0}^{\infty} F(x) G(x) dx = \int_{0}^{\infty} f_{\mu_{r}}(y^{-1}) g_{\mu_{r}}(y^{-1}) y^{-2} dy$$

$$= \int_{0}^{\infty} f_{\mu_{r}}(y) g_{\mu_{r}}(y) dy, \quad r = 1, 2, ... n.$$

All the conditions of De la Vallée Poussin's theorem (see Carslaw, H.S. Introduction to the Theory of Fourier's Series and Integrals. Art. 89, p.209) are satisfied and a change in the order of integrations can be effected.

COROLLARY 1. Let

$$F(x) = \int_{0}^{\infty} f(t) \tilde{\omega}_{\mu_{1}, \dots, \mu_{n}}(xt) dt, G(x) = \int_{0}^{\infty} g(t) \tilde{\omega}_{\mu_{1}, \dots, \mu_{n}}(xt) dt.$$

Then

$$\int_{0}^{\infty} F(x)G(x)dx = \int_{0}^{\infty} f_{\mu_{1}}, \dots \mu_{r}(y) g_{\mu_{1}}, \dots \mu_{r}(y)dy, r = 1, 2, \dots n,$$

$$R(\mu_{r}) > -1, r = 1, 2, \dots r,$$

under the conditions of the theorem.

COROLLARY 2. Let

$$F(x) = \int_{0}^{\infty} f(t) \frac{\partial}{\partial \mu_{1}}, \dots \mu_{m} (xt) dt$$

$$G(x) = \int_{0}^{\infty} g(t) \frac{\partial}{\partial \mu_{1}}, \dots \mu_{n} (xt) dt, \quad m < n,$$

then

$$\int_{0}^{\infty} F(x)G(x)dx = \int_{0}^{\infty} f_{\mu_{1},\mu_{2},\dots,\mu_{r}}(y)g_{\mu_{1},\dots,\mu_{r},\mu_{m+1},\dots,\mu_{n}}(y)dy,$$

$$r < m < n, \qquad R(\mu_{s}) > -1, \qquad s = 1, 2, \dots, n,$$

under the conditions of the theorem.

Proof.

$$\int_{0}^{\infty} F(x)G(x)dx = \int_{0}^{\infty} G(x)dx \int_{0}^{\infty} f_{\mu_{1}, \dots \mu_{r}}(y)\widetilde{\omega}_{\mu_{r+1}, \dots \mu_{m}}(x/y)y^{-1}dy$$

$$= \int_{0}^{\infty} f_{\mu_{1}, \dots \mu_{r}}(y)y^{-1}dy \int_{0}^{\infty} g(t)\widetilde{\omega}_{\mu_{1}, \dots \mu_{n}}^{\mu_{r+1}, \dots \mu_{m}}(yt) ydt$$

$$= \int_{0}^{\infty} f_{\mu_{1}, \dots \mu_{r}}(y) g_{\mu_{1}, \dots \mu_{r}, \mu_{m+1}, \dots \mu_{n}}(y)dy.$$

## Examples

1. Let f(t) = g(t), so that F(x) = G(x).

$$\int_{0}^{\infty} x^{-\frac{1}{4}} \frac{1}{(1+ax)^{2}} \sum_{\omega_{2b-\frac{1}{2},0}}^{\infty} (p^{\frac{1}{2}} x^{\frac{1}{2}}) dx$$

\* The second integral on the right can be deduced from equation (2).

$$= \frac{4}{a} p^{\frac{1}{4}} K_{2b - \frac{1}{2}} (2^{\frac{1}{2}} a^{-\frac{1}{4}} \frac{1}{p^{4}}) J_{2b - \frac{1}{2}} (2^{\frac{1}{2}} a^{-\frac{1}{4}} \frac{1}{p^{4}}),$$

or

$$\int_{0}^{\infty} \frac{1}{x^{2}} (1 + a^{2}x^{2})^{-\frac{3}{2}} \frac{1}{2} \frac{1}{2} \int_{0}^{\infty} (px) dx$$

$$= 2a^{-2}\frac{\frac{1}{2}}{p}K \left(\sqrt{\frac{2p}{a}}\right) J \left(\sqrt{\frac{2p}{a}}\right) .$$

$$f(x) = \frac{\frac{1}{2}}{x^2/(1+a^2x^2)^2}$$
,  $F(x) = 2a^{-2}x^{\frac{1}{2}}K$   
 $2b - \frac{1}{2}(\sqrt{\frac{2x}{a}}) J$   
 $2b - \frac{1}{2}(\sqrt{\frac{2x}{a}})$ 

and

$$\int_{0}^{\infty} x \left\{ K_{2b - \frac{1}{2}} \left( \sqrt{\frac{2x}{a}} \right) J_{2b - \frac{1}{2}} \left( \sqrt{\frac{2x}{a}} \right) \right\}^{2} dx = \frac{a^{4}}{4} \int_{0}^{\infty} \frac{x dx}{(1 + a^{2}x^{2})^{3}} = \frac{a^{2}}{16},$$

$$R(b) > -\frac{1}{4}.$$

2. Let 
$$f(x) = F(x) = x^{\frac{1}{2}} K_{v}(x)$$
 which is  $R_{-v, v}$ ,  $-1 < R(v) < 1$ .

$$f_{v}(y) = \int_{0}^{\infty} x^{\frac{1}{2}} K_{v}(x) \sqrt{xy} J_{v}(xy) dx = y^{v+\frac{1}{2}}/(1+y^{2}).$$

\* See [1], References. This can be proved by an application of Parseval's theorem on Mellin Transform.

Therefore

$$\int_{0}^{\infty} \{f_{v}(y)\}^{2} dy = \int_{0}^{\infty} \frac{y^{2v+1}}{(1+y^{2})^{2}} dy = \frac{v\pi}{2 \sin v\pi}.$$

Also

$$\int_{0}^{\infty} \{F(x)\}^{2} dx = \int_{0}^{\infty} \{x^{2} K_{v}(x)\}^{2} dx = \frac{v\pi}{2 \sin v\pi}.$$

Hence

$$\int_{0}^{\infty} \left\{ F(x) \right\}^{2} dx = \int_{0}^{\infty} \left\{ f(y) \right\}^{2} dy.$$

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