

**UNIQUENESS OF SOLUTIONS OF IMPROPERLY
POSED PROBLEMS FOR SINGULAR
ULTRAHYPERBOLIC EQUATIONS**

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1. Introduction

In [1], Owen gave sufficient conditions for the uniqueness of certain mixed problems having elliptic and hyperbolic nature for the ultrahyperbolic equation. Recently, Diaz and Young [2] has obtained necessary and sufficient conditions for the uniqueness of solutions of the Dirichlet and Neumann problems involving the more general ultrahyperbolic equation

$$\Delta u - D_j(a_{jk}D_k u) + cu = 0$$

The purpose of this paper is to present corresponding uniqueness conditions for the Dirichlet and Neumann problems for the singular ultrahyperbolic equation

$$(1) \quad Lu \equiv u_{tt} + (\alpha/t)u_t + \Delta u - D_j(a_{jk}D_k u) + cu = 0$$

for all values of the parameter α , $-\infty < \alpha < \infty$. The symbol Δ denotes the Laplace operator in the variables x_1, \dots, x_m , D_j indicates partial differentiation with respect to the variable y_j ($1 \leq j \leq n$), and the summation convention is adopted for repeated indices including $(\partial_i u)^2$, where ∂_i denotes differentiation with respect to the variable x_i .

The boundary value problems will be considered in the domain $Q = X^* \times Y$ where X^* is the parallelepiped defined by $0 < t < T$, $0 < x_i < a_i$ ($1 \leq i \leq m$), and Y is a bounded domain in the space y_1, \dots, y_n . The parallelepiped defined by $0 < x_i < a_i$ ($1 \leq i \leq m$) will be denoted by X . For brevity, we write $x = (x_1, \dots, x_m)$, $y = (y_1, \dots, y_n)$, and denote a point in Q by (t, x, y) .

Throughout this paper, we assume that the coefficients a_{jk} and c depend only on the variables y_1, \dots, y_n , and are continuous functions of these variables with

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$c \geq 0$ in Y . As in [2] and [4], we also assume that the matrix (a_{jk}) is symmetric, positive definite, and that a_{jk} , c and the domain Y are sufficiently regular in order to allow the application of the divergence theorem and to ensure the existence of a complete set of eigenfunctions of class $C^2(Y) \cap C^1(\bar{Y})$ for the eigenvalue problems that will be needed below. By a solution of a boundary value problem considered here we shall mean a function $u \in C^2(Q) \cap C^1(\bar{Q})$ which satisfies the differential equation and the boundary condition of the problem.

The results given here include not only those obtained in [2], but also those derived by Dunninger and Zachmanoglov [3], [4], Sigillito [5] and Young [6] in the case of the normal hyperbolic equation.

2. The Dirichlet problem

We consider first the homogeneous Dirichlet problem

$$(2) \quad Lu = 0 \text{ in } Q, \quad u = 0 \text{ on } \partial Q$$

Corresponding to various ranges of the parameter α , we shall prove uniqueness of solution by showing that every solution of the problem vanishes identically in Q . We begin by stating a lemma which characterizes every smooth solution of the equation (1) for $\alpha \neq 0$.

LEMMA. *If $\alpha \neq 0$, then every solution u of (1) belonging to C^2 for $t > 0$ and to C^1 for $t \geq 0$ satisfies the condition $u_t(0, x, y) = 0$.*

This lemma can be proved by following, almost step for step, the method employed by Fox [7] in establishing the same property for the corresponding singular normal hyperbolic equation in the case that (a_{jk}) is the identity matrix, using the domain Q .

THEOREM 1. *If $\alpha > 0$, then every solution of the problem (2) vanishes identically in Q .*

PROOF. Let u be a solution of the problem (2). We integrate the identity

$$\begin{aligned} 0 = 2u_t Lu &= [u_t^2 - (\partial_i u)^2 + a_{jk} D_j u D_k u + cu^2]_t \\ &\quad + 2\partial_i(u_t \partial_i u) - 2D_j(a_{jk} u_t D_k u) \\ &\quad + (2\alpha/t)u_t^2 \end{aligned}$$

over Q and use the divergence theorem to obtain

$$(3) \quad \begin{aligned} &\int_{X \times Y} (u_t^2 - (\partial_i u)^2 + a_{jk} D_j u D_k u + cu^2) \Big|_{t=0}^{t=T} dx dy \\ &\quad + \int_{\partial Q} (2u_t \partial_i u v_i - 2u_t a_{jk} D_k u v_j^*) dS \end{aligned}$$

$$+ 2\alpha \int_Q (u_t^2/t) dt dx dy = 0$$

where v_i and v_j^* denote the components of the outward normal vector on ∂X and ∂Y , respectively. By the lemma and the fact that u vanishes on ∂Q , (3) reduces to

$$(4) \quad \int_{X \times Y} u_t^2(T, x, y) dx dy + 2\alpha \int_Q (u_t^2/t) dt dx dy = 0$$

Since $\alpha > 0$, this implies that $u_t \equiv 0$ in Q , that is, u is independent of t . But $u = 0$ on $t = 0$, hence $u \equiv 0$ in Q .

THEOREM 2. *Let λ_r ($r = 1, 2, \dots$) be the eigenvalues of the problem*

$$D_j(a_{jk} D_k v) - cv + \lambda v = 0 \text{ in } Y$$

(5)

$$v = 0 \text{ on } \partial Y.$$

If $\alpha \leq 0$, then every solution of the problem (2) vanishes identically in Q if and only if

$$(6) \quad J_{(1-\alpha)/2}(\mu^{\frac{1}{2}} T) \neq 0$$

for any real number $\mu \neq 0$ and nonzero integers p_1, \dots, p_m such that

$$(7) \quad \mu + \sum_{i=1}^m (p_i \pi / a_i)^2 = \lambda_r,$$

where $J_\alpha(t)$ is the Bessel's function of the first kind of order α .

PROOF. Suppose that there exist an eigenvalue λ_s of (5), a real number $\mu_s \neq 0$, and nonzero integers q_1, \dots, q_m satisfying (7) such that

$$(8) \quad J_{(1-\alpha)/2}(\mu_s^{\frac{1}{2}} T) = 0.$$

Let v_s be an eigenfunction of (5) corresponding to λ_s , and define

$$(9) \quad \phi(x; q) = \prod_{i=1}^m \sin(q_i \pi x_i / a_i).$$

Then it is readily verified that the function

$$u(t, x, y) = t^{(1-\alpha)/2} J_{(1-\alpha)/2}(\mu_s^{\frac{1}{2}} t) \phi(x; q) v_s(y)$$

is a nontrivial solution of the problem (2).

Conversely, suppose that the conditions (6) and (7) hold. Let us integrate the identity

$$\begin{aligned}
 wLu - uMw &= (wu_t - w_tu + \alpha uw/t)_t \\
 &\quad + \partial_i(w\partial_iu - u\partial_iw) \\
 &\quad - D_j[a_{jk}(wD_ku - uD_kw)]
 \end{aligned}$$

over $Q_s = X_s^* \times Y$, where X_s^* is the parallelepiped defined by $0 < s \leq t < T$, $0 < x_i < a_i$ ($1 \leq i \leq m$), and M is the adjoint operator of L given by

$$\begin{aligned}
 Mw &= w_{tt} - \alpha(w/t)_t + \Delta w \\
 &\quad - D_j(a_{jk}D_kw) + cw.
 \end{aligned}$$

By the divergence theorem, we have

$$\begin{aligned}
 (10) \quad &\int_{Q_s} [wLu - uMw] dt dx dy \\
 &= \int_{\partial Q_s} [(wu_t - w_tu + \alpha uw/t)v_t + (w\partial_iu - u\partial_iw)v_i \\
 &\quad - a_{jk}(wD_ku - uD_kw)v_j^*] dS.
 \end{aligned}$$

Now let u be a solution of (2) and for any choice of λ_r , $u \neq 0$, and nonzero integers p_1, \dots, p_m satisfying (6) and (7), let

$$w(t, x, y) = t^{(1+\alpha)/2} J_{(1-\alpha)/2}(\mu^{\frac{1}{2}}t) \phi(x; p) v_r(y)$$

where ϕ is defined in (9) and v_r is an eigenfunction associated with λ_r . Since $Lu = 0$ and

$$Mw = -t^{(1+\alpha)/2} J_{(1-\alpha)/2}(\mu^{\frac{1}{2}}t) \phi(x; p) [D_j(a_{jk}D_kv_r) - cv_r + \lambda_r v_r] = 0,$$

the left hand side of (10) vanishes. Moreover, since $u = 0$ on ∂Q and $w = 0$ on $X^* \times \partial Y$ and $\partial X \times Y$, equation (10) becomes

$$(11) \quad \int_{X \times Y} (wu_t - w_tu + \alpha uw/t) \Big|_{t=s}^{t=T} dx dy = 0.$$

We now let s approach zero. Since both w_t and w/t are bounded at $t = 0$, and u vanishes there, we obtain in the limit

$$(12) \quad T^{(1+\alpha)/2} J_{(1-\alpha)/2}(\mu^{\frac{1}{2}}T) \int_{X \times Y} u_t(T, x, y) \phi(x; p) v_r(y) dx dy = 0.$$

In view of (6) and the completeness of the sets of eigenfunctions $\{\prod_{i=1}^m \sin(p_i \pi x_i / a_i)\}$ and $\{v_r\}$ in X and Y , respectively, (12) implies that $u_t(T, x, y) = 0$. With this additional information, we can now show that $u \equiv 0$ in Q .

Let us integrate the identity

$$(13) \quad 0 = (2tu_t + u)Lu = [t(u_t^2 - (\partial_i u)^2 + a_{jk}D_j u D_k u + cu^2) + u(u_t + \frac{1}{2}\alpha u/t), \\ + \partial_i[(2tu_t + u)\partial_i u] - D_j[a_{jk}(2tu_t + u)D_k u] + 2(\alpha - 1)u_t^2 + \frac{1}{2}\alpha u^2/t^2$$

over Q_s and pass to the limit as $s \rightarrow 0$. Since $u = 0$ on ∂Q , $u_t = 0$ on $t = 0$ and $t = T$, all surface integrals arising from the integration vanish in the limit, so that we are left with the convergent integral

$$\int_Q [2(\alpha - 1)u_t^2 + \frac{1}{2}\alpha u^2/t^2] dt dx dy = 0.$$

Since $\alpha \leq 0$, this yields the result that $u \equiv 0$ in Q .

3. The Neumann problem

We consider next the homogeneous Neumann problem

$$(14) \quad Lu = 0 \text{ in } Q, \quad \partial u / \partial n = 0 \text{ on } \partial Q,$$

where $\partial u / \partial n$ denotes the conormal derivative

$$\partial u / \partial n = a_{jk} D_k u v_j^*$$

on the part $X^* \times \partial Y$ of ∂Q .

THEOREM 3. *Let λ_r ($r = 1, 2, \dots$) be the nonzero eigenvalues of the problem*

$$(15) \quad D_j(a_{jk} D_k v) - cv + \lambda v = 0 \text{ in } Y, \\ \partial v / \partial n = 0 \text{ on } \partial Y.$$

Then every solution u of the problem (14) vanishes identically (or $u = \text{const.}$ if $c \equiv 0$) for $\alpha \geq 0$ if and only if

$$(16) \quad J_{(1+\alpha)/2}(\mu^{\frac{1}{2}} T) \neq 0$$

for any real number $\mu \neq 0$ and integers p_1, \dots, p_m satisfying (7).

PROOF. The condition (16) is actually necessary for any value of the parameter α . In fact, if there exist a nonzero eigenvalue λ_s of (15), a real number $\mu_s \neq 0$, and integers q_1, \dots, q_m satisfying (7) such that

$$(17) \quad J_{(1+\alpha)/2}(\mu_s^{\frac{1}{2}} T) = 0$$

then the function

$$(18) \quad u(t, x, y) = t^{(1-\alpha)/2} J_{(\alpha-1)/2}(\mu_s^{\frac{1}{2}} t) \psi(x; q) v_s(y)$$

constitutes a nontrivial solution of the problem (14) for any value of α . Here

$$(19) \quad \psi(x; q) = \prod_{i=1}^m \cos(q_i \pi x_i / a_i)$$

and v_s is an eigenfunction corresponding to λ_s . Indeed, by (15) it is easily shown that (18) satisfies $Lu = 0$, $\partial u / \partial n = 0$ on $X^* \times \partial Y$, and $\partial_i u = 0$ on $x_i = 0$, $x_i = a$ ($1 \leq i \leq m$). Moreover, since

$$u_t = -\mu_s^{\frac{1}{2}} t^{(1-\alpha)/2} J_{(1+\alpha)/2}(\mu_s^{\frac{1}{2}} t) \psi(x; q) v_s(y) = O(t),$$

it follows that $u_t(0, x, y) = 0$ and by (17) $u_t(T, x, y) = 0$. Thus (18) is a nontrivial solution of the problem (14).

On the other hand, let $\alpha \geq 0$ and assume that the condition (16) holds. Let λ_r be a nonzero eigenvalue of (15) with the corresponding eigenfunction v_r . For any choice of real number $\mu \neq 0$ and integers p_1, \dots, p_m satisfying (7) and (16), let

$$(20) \quad w(t, x, y) = t^{(1+\alpha)/2} J_{(\alpha-1)/2}(\mu^{\frac{1}{2}} t) \psi(x; p) v_r(y)$$

where ψ is given by (19). By direct differentiation, it is readily verified that $Mw = 0$ in Q , $\partial w / \partial n = 0$ on ∂Q except on $t = 0$ and $t = T$. Hence, if u is a solution of (14), substitution of (20) for w in (10) leads again to the integral (11). Since

$$w_t = [\alpha t^{(\alpha-1)/2} J_{(\alpha-1)/2}(\mu^{\frac{1}{2}} t) - \mu^{\frac{1}{2}} t^{(1+\alpha)/2} J_{(\alpha+1)/2}(\mu^{\frac{1}{2}} t)] \psi(x; p) v_r(y)$$

it follows that

$$\begin{aligned} -w_t + \alpha w / t &= \mu^{\frac{1}{2}} t^{(1+\alpha)/2} J_{(\alpha+1)/2}(\mu^{\frac{1}{2}} t) \psi(x; p) v_r(y) \\ &= O(t^{\alpha+1}). \end{aligned}$$

Therefore, as s is allowed to approach zero in (11), we obtain in the limit

$$\mu^{\frac{1}{2}} T^{(1+\alpha)/2} J_{(\alpha+1)/2}(\mu^{\frac{1}{2}} T) \int_{X \times Y} u(T, x, y) \psi(x; p) v_r(y) dx dy = 0,$$

By the hypothesis (16) and the completeness of the sets of eigenfunctions $\{\prod_{i=1}^m \cos(p_i \pi x_i / a_i)\}$ and $\{v_r\}$ in X and Y , respectively, we conclude that $u(T, x, y) = 0$ if $c > 0$ and $u(T, x, y) = \text{const.}$ if $c \equiv 0$. Notice that in the case $c \equiv 0$, the problem (15) has the eigenfunction $v = 1$ corresponding to the eigenvalue $\lambda = 0$.

Let us consider the case $c > 0$. It remains to be shown that $u \equiv 0$ in Q . For this purpose, we note that the identity (13) no longer applies. We integrate instead the identity

$$\begin{aligned} [2t^{\alpha+1} u_t + (\alpha + 1)t^\alpha u] Lu &= [t^{\alpha+1}(u_t^2 - (\partial_i u)^2 + a_{jk} D_j u D_k u + cu^2) \\ &\quad + (\alpha + 1)t^\alpha uu_t]_t + \partial_i [t^\alpha (2tu_t + (\alpha + 1)u) \partial_i u] \end{aligned}$$

$$-D_j[t^\alpha a_{jk}(2tu_t + (\alpha + 1)u)D_k u] - 2t^\alpha u_t^2$$

over Q and apply the divergence theorem. Because u is a solution of (14) and $u = 0$ at $t = T$, it is clear that all surface integrals arising from the integration vanish. Thus we have

$$-2 \int_Q t^\alpha u_t^2 dt dx dy = 0$$

from which the conclusion that $u \equiv 0$ in Q follows.

If $c \equiv 0$, then the above argument gives $u = \text{const.}$ in Q .

4. Concluding remarks

By using the same technique, it is possible to prove uniqueness theorems for equation (1) subject to mixed boundary conditions of the type considered in [2] with respect to the variables x , y and with either the condition $u = 0$ or $u_t = 0$ on $t = 0$ and $t = T$.

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